

University of Illinois at Urbana-Champaign, June 13-23, 2005

Introduction to the SIESTA method

Some technicalities

Linear-scaling DFT based on Numerical Atomic Orbitals (NAOs)

- Born-Oppenheimer
- DFT
- Pseudopotentials
- Numerical atomic orbitals
- Numerical evaluation of matrix elements

- relaxations, MD, phonons.
- LDA, GGA
- norm conserving, factorised.
- finite range

P. Ordejon, E. Artacho & J. M. Soler , Phys. Rev. B 53, R10441 (1996)J. M.Soler et al, J. Phys.: Condens. Matter 14, 2745 (2002)

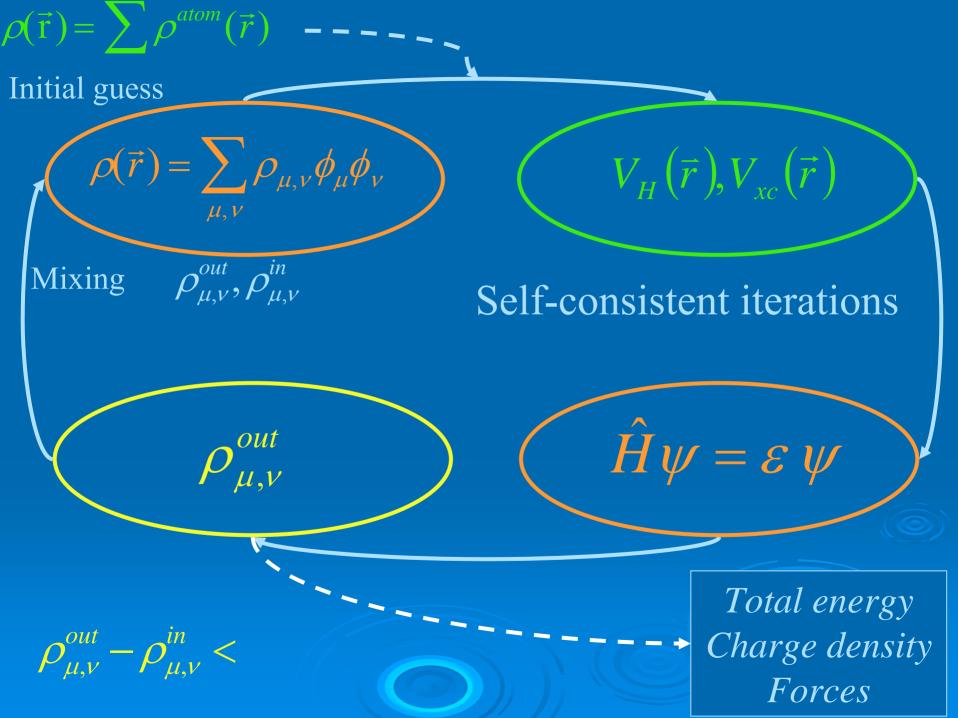
Matrix equations:

$$\psi_{n}(\vec{r}) \approx \sum_{\mu} c_{\mu n} \phi_{\mu}(\vec{r}) \qquad \text{Expand in terms of a finite set of known wave-functions} \qquad \phi_{\mu}(\vec{r})$$

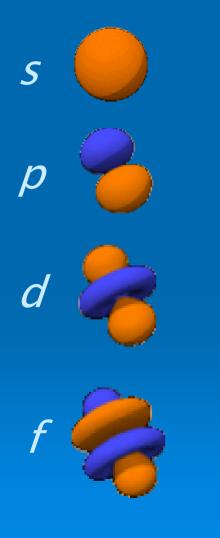
$$\psi_{n}(\vec{r}) = \varepsilon_{n} \psi_{n}(\vec{r}) \longrightarrow \sum_{\mu} c_{\mu n} \hat{h} \phi_{\mu}(\vec{r}) = \varepsilon_{n} \sum_{\mu} c_{\mu n} \phi_{\mu}(\vec{r})$$

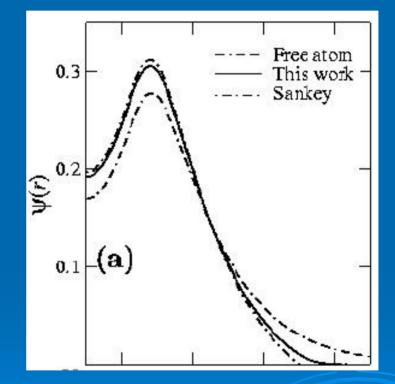
$$\text{Def. } h_{\nu \mu} = \int \phi_{\nu}^{*}(\vec{r}) \hat{h} \phi_{\mu}(\vec{r}) d^{3}\vec{r} \text{ and } S_{\nu \mu} = \int \phi_{\nu}^{*}(\vec{r}) \phi_{\mu}(\vec{r}) d^{3}\vec{r}$$

$$\sum_{\mu} h_{\nu \mu} c_{\mu n} = \varepsilon_{n} \sum_{\mu} S_{\nu \mu} c_{\mu n}$$



Basis Set: Atomic Orbitals





Strictly localised (zero beyond a cut-off radius)

Long-range potentials $H = T + V_{ion}(r) + V_{nl} + V_{H}(r) + V_{xc}(r)$

_ong range

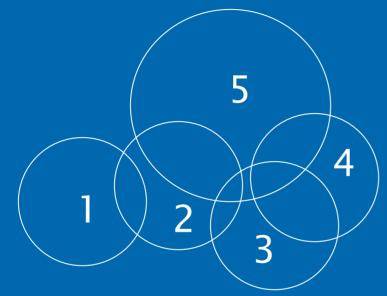
$$\begin{split} V_{na}(\mathbf{r}) &= V_{ion}(\mathbf{r}) + V_{H}[\rho_{atoms}(\mathbf{r})] & \text{Neutral-atom potential} \\ \delta V_{H}(\mathbf{r}) &= V_{H}[\rho_{SCF}(\mathbf{r})] - V_{H}[\rho_{atoms}(\mathbf{r})] \end{split}$$

 $\mathbf{H} = \mathbf{T} + \mathbf{V}_{nl} + \mathbf{V}_{na}(\mathbf{r}) + \delta \mathbf{V}_{H}(\mathbf{r}) + \mathbf{V}_{xc}(\mathbf{r})$

Two-center integrals

Grid integrals

Sparsity



1 with 1 and 2 2 with 1,2,3, and 5 3 with 2,3,4, and 5 4 with 3,4 and 5 5 with 2,3,4, and 5

Non-overlap interactions Basis orbitals

 $S_{\mu\nu}$ and $\,H_{\mu\nu}$ are sparse

 $\rho_{\mu\nu}$ is not strictly sparse but only a sparse subset is needed

KB pseudopotential projector

> Two-center integrals Convolution theorem $S(\mathbf{R}) \equiv \left\langle \phi_1 \, \middle| \, \phi_2 \right\rangle = \int \phi_1(\mathbf{r}) \, \phi_2(\mathbf{r} - \mathbf{R}) \, d\mathbf{r}$ $\phi(\mathbf{k}) = \frac{1}{(2\pi)^{2/3}} \int \phi(\mathbf{r}) e^{-i\mathbf{k}\mathbf{r}} d\mathbf{r}$ $S(\mathbf{R}) = \int \phi_1(\mathbf{k}) \, \phi_2(\mathbf{k}) \, e^{i\mathbf{k}\mathbf{R}} d\mathbf{k}$

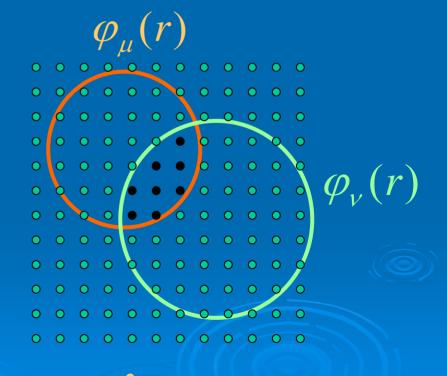
Grid work

$$\psi_{i}(\mathbf{r}) = \sum_{\mu} c_{i\mu} \phi_{\mu}(\mathbf{r})$$

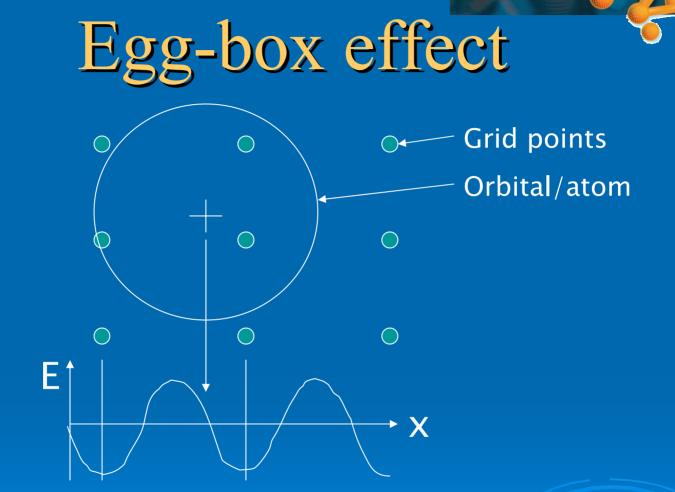
$$\rho_{\mu\nu} = \sum_{i} c_{i\mu} c_{i\nu}$$

$$\rho(\mathbf{r}) = \sum_{i} \psi_{i}^{2}(\mathbf{r}) = \sum_{\mu\nu} \rho_{\mu\nu} \phi_{\mu}(\mathbf{r}) \phi_{\nu}(\mathbf{r})$$

 $\rho(\mathbf{r}) \rightarrow V_{xc}(\mathbf{r})$ $\delta\rho(\mathbf{r}) = \rho_{SCF}(\mathbf{r}) - \rho_{atoms}(\mathbf{r})$ $\delta\rho(\mathbf{r}) \xrightarrow{FFT} \delta V_{H}(\mathbf{r})$

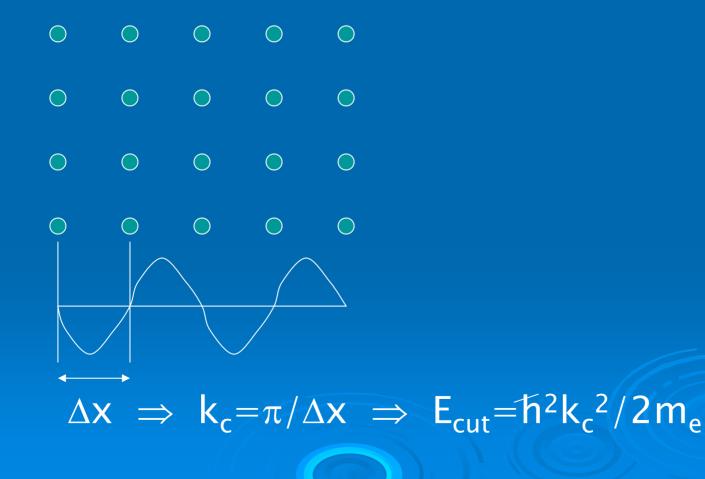


 $H_{\mu\nu} = \left\langle \varphi_{\mu}(r) | V | \varphi_{\nu}(r) \right\rangle = \int \varphi_{\mu}(r) V(r) \varphi_{\nu}(r)$

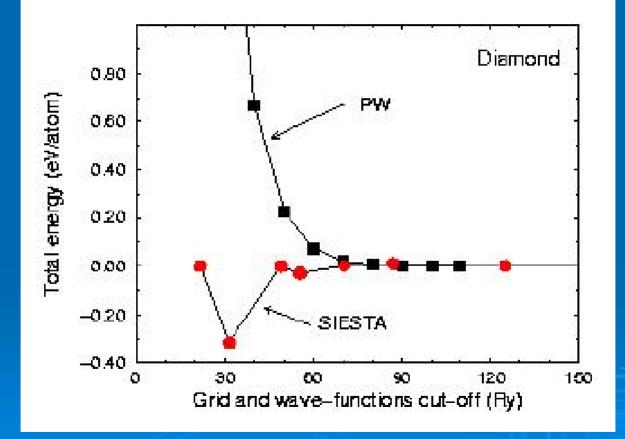


Affects more to forces than to energyGrid-cell sampling

Grid smoothness: energy cutoff



Grid smoothness: convergence



 $E_{cut} = (\pi / \Delta x)^2$

Boundary conditions

Isolated object (atom, molecule, cluster):
 Open boundary conditions (defined at infinity)

> 3D Periodic object (crystal):
 Periodic Boundary Conditions

> Mixed:

- 1D periodic (chains)
- 2D periodic (slabs)

