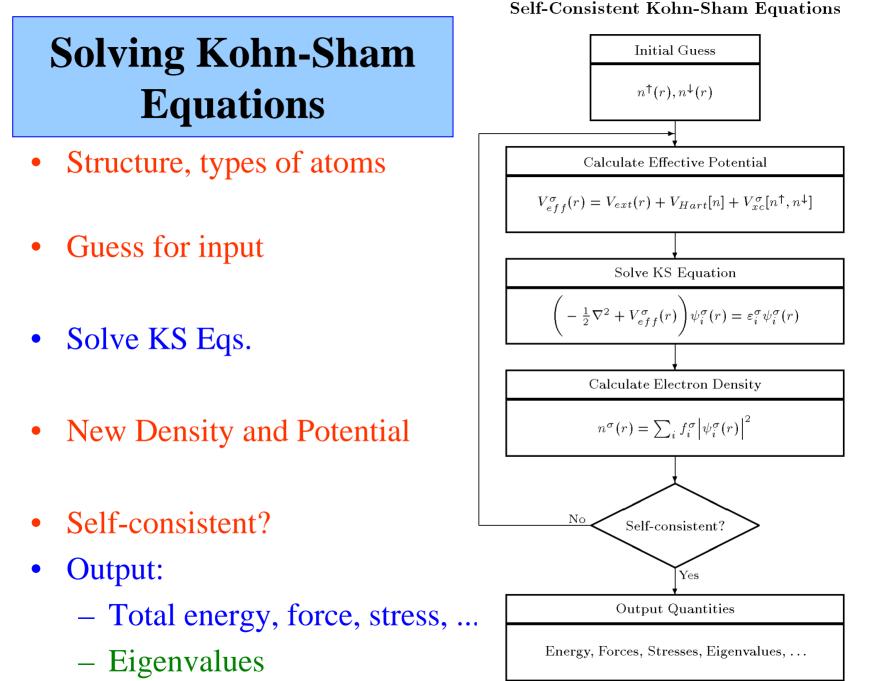
Electron bands in crystals Pseudopotentials, Plane Waves, Local Orbitals

Richard M. Martin – UIUC

Lecture at Summer School Hands-on introduction to Electronic Structure Materials Computation Center University of Illinois – June, 2005

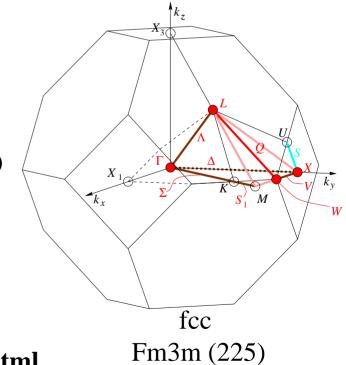
Crystals and electron bands Pseudopotentials Plane Wave Methods Local Orbital Methods Reference: Electronic Structure: Basic Theory and Practical Methods,

Richard M. Martin (Cambridge University Press, 2004)



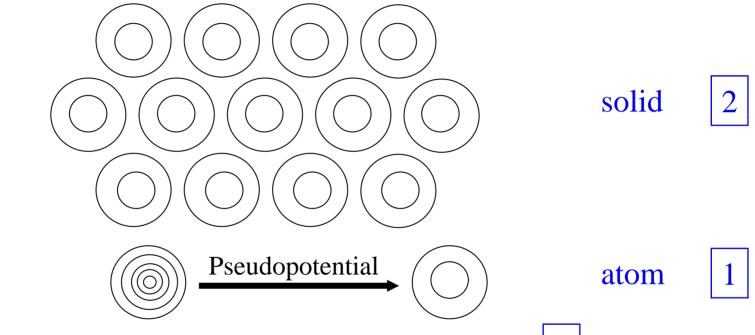
Structure

- "External Potential" for electrons determined by structure
 - Lattice periodicty
 - Positions of atoms in primitive cell
 - Chemical identity of each atom atomic number
- Symmetry
 - Translation symmetry
 - Lattice Reciprocal lattice
 - Bloch Theorem
 - States labeled by k in Brillouin Zone (BZ)
- See Notes and examples in Lab
 - Example fcc crystals
- Useful Web site
 - BZ with labels for all 230 space groups http://www.cryst.ehu.es/cryst/get_kvec.html



Pseudopotentials

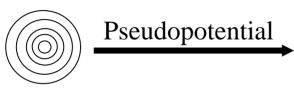
• Pseudopotential Method – replace each potential



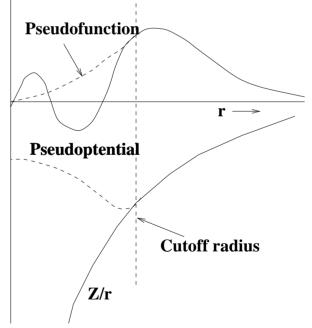
- 1 Generate Pseudopotential in atom (spherical) 2 use in solid
- Pseudopotential can be constructed to be weak
 - Can be chosen to be smooth
 - Solve Kohn-Sham equations in solid directly in Fourier space

Pseudopotentials

- Norm-Conserving Pseudopotential (NCPP)
 - Hamann, Schluter, Chaing



- Generate weak pseudopotential in atom with same scattering properties for valence states as the strong all-electron potential
- Conditions
 - Potential same for $r > R_c$
 - Pseudofunction "norm-conserving" for $r < R_c$
- Codes available for generating potentials



Pseudopotentials

• Web Site: http://www.tddft.org/programs/octopus/pseudo.php



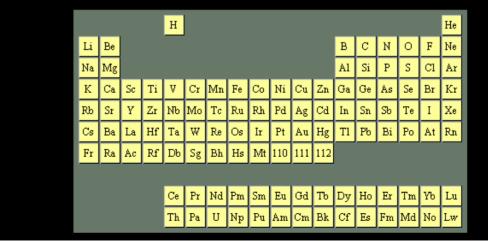
News | Download | Developers | Pseudopotentials | Contributors | Info

Pseudopotentials

Octopus uses non-local, norm-conserving Troullier-Martins pseudopotentials. You can generate them with **José Luis Martins**' atomic **program**, or use the pseudopotentials distributed from this web page.

Obviously you will use the pseudopotential at your own risk, and we do not feel responsible for any (bad) result that could outcome from the use of these files. I suppose that you are grown-up and understand that ;))

Just click on an element



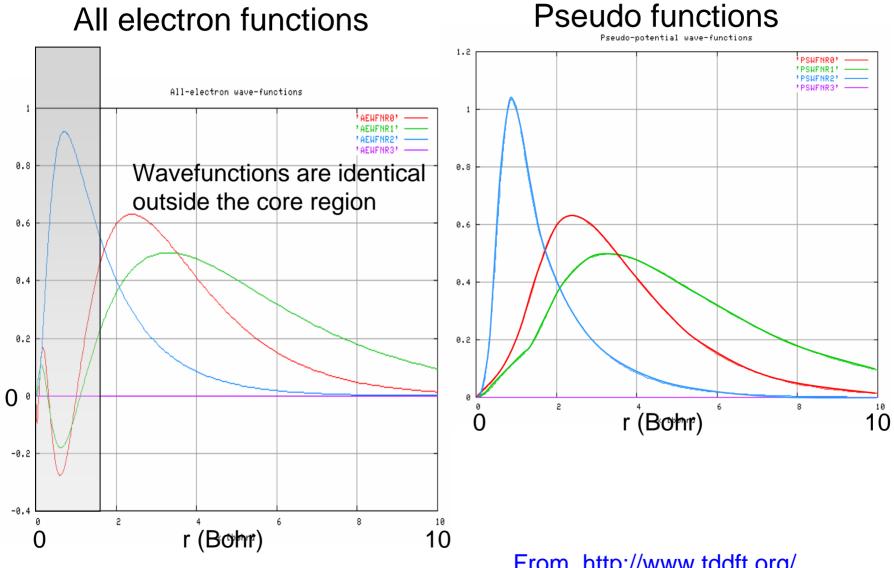
Not everything works, but part does!

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Use with care!

Always Test!

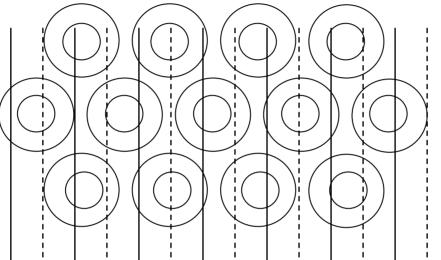
Fe – s,p,d valence wavefunctions – r $\Psi(r)$



From http://www.tddft.org/

Plane Waves

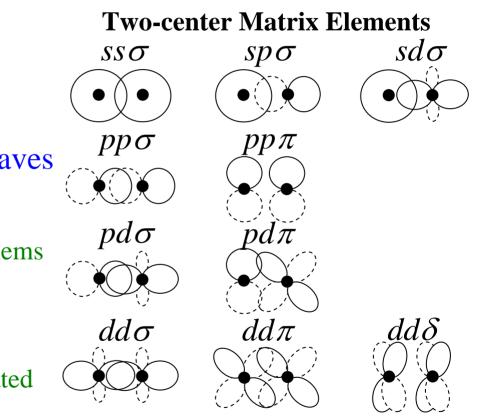
• Pseudopotential method with plane waves



- See Notes and Lab for equations and descriptions TBPW
 - Examples with approximate empirical pseudopotentials
- Full self-consistent calculations ABINIT, other codes
 - Method of choice for many reasons Simplicity and the speed of Fast Fourier Transforms
 - Forces, phonons, response functions, ...
 - Simulations solids, liquids,

Localized orbitals

- Intuitive appeal of atomic-like states
 - Gaussian basis widely used in chemistry
 - Numerical radial orbitals used in SIESTA
- More efficient than plane waves
 - Very useful for complicated problems
 - Localized, non-periodic problems
- Tight-binding
 - Simple understanding
 - Interpolation of bands calculated by other methods
 - Very fast and applicable to very large systems



Localized Orbitals & Tight binding

• Localized orbitals

- Full methods Calculate all quantities done in SIESTA
- Tight binding Take matrix elements as parameters
- Very useful qualitative quantitative

Define localized orbitals $\chi_m^*(\mathbf{r} - \tau_m)$ centered at site τ_m), where m = 1, 2, ... labels all the orbitals. Then

$$H_{m,m'} = \int d\mathbf{r} \ \chi_m^*(\mathbf{r} - \tau_m) \widehat{H} \chi_{m'}(\mathbf{r} - \tau_{m'}), \ (1)$$

and the overlap matrix is given by

$$S_{m,m'} = \int d\mathbf{r} \ \chi_m^* (\mathbf{r} - \tau_m) \chi_{m'} (\mathbf{r} - \tau_{m'}). \quad (2)$$

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Localized Orbitals & Tight binding

In a crystal one can define basis states with wavevector ${\bf k},$

$$\chi_{m\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{T}} e^{i\mathbf{k}\cdot\mathbf{T}} \chi_m(\mathbf{r} - (\tau_m + \mathbf{T})). \quad (1)$$

The solution can be written

$$\psi_{i\mathbf{k}}(\mathbf{r}) = \sum_{m} c_{m}(\mathbf{k}) \chi_{m\mathbf{k}}(\mathbf{r}),$$
 (2)

and the secular equation for wavevector ${\bf k}$ is

$$\sum_{m'} \left[H_{m,m'}(\mathbf{k}) - \varepsilon_i(\mathbf{k}) S_{m,m'}(\mathbf{k}) \right] c_{i,m'}(\mathbf{k}) = 0.$$
(3)
See Notes and Lab for examples – TBPW code

Total energy and force in tight binding

The total energy can always be written

$$E_{tot} = \sum_{i} \varepsilon_{i} f(\varepsilon_{i}) + F[n].$$
 (1)

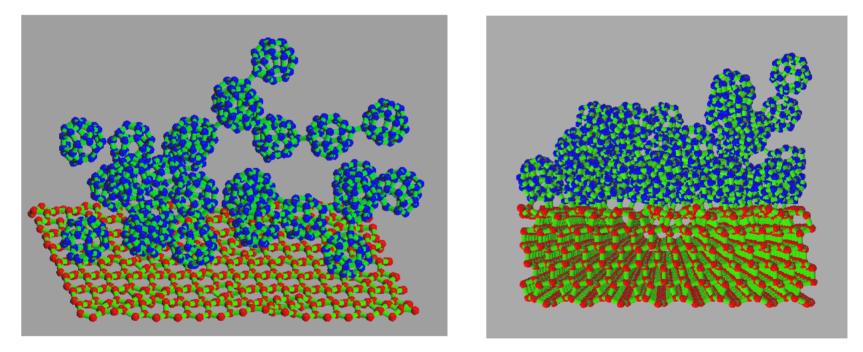
where ε_i are eigenvalues, f is the Fermi function, and F[n] is the replusive sum of nuclei intercations and double counting corrections. F[n] can be approximated as 2-body terms,

$$E_{tot} = \sum_{i} \varepsilon_{i} f(\varepsilon_{i}) + \sum_{I < J} f(|\mathbf{R}_{I} - \mathbf{R}_{J}|).$$
(2)

- Key Points
 - Adjustable parameters in hamiltonian, overlap and f terms
 - Can vary from "quick and dirty" to quite accurate
 - Very fast calculations bands, energy, forces,

Deposition of C₂₈ Buckyballs on Diamond

Example using tight-binding and "Order-N" methods



 Simulations with ~ 5000 atoms, TB Hamiltonian from Xu, et al. (A. Canning, G.~Galli and J.Kim, Phys.Rev.Lett. 78, 4442 (1997).

Conclusions

- Solution of the independent-particle Kohn-Sham equations is the essential step in DFT calculations
 - Most-time-consuming part
- All DFT calculations share many common features
 - Stucture of nuclei, must specify external potential
- Differences
 - Basis Plane wave, localized, ...
 - All-electron, pseudopotential
 - Diagonalization of matrix vs iterative solution
- Many features can be illustrated by simple examples and a pedagogical code TBPW
 - Lab with TBPW
- Full calculations illustrated by ABINIT, SIESTA