

A first-principles approach to orbital accumulation and orbital transport

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UNIVERSITET

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Outline

Aim: Develop electronic structure theory to compute and predict unusual spin and orbital effects

- Motivation: Spin-orbit torques
- Electronic structure methodology
- Results: Orbital Hall effect & orbital Nernst effect in 40 elements
- Magnetic spin Hall effect and magnetic orbital Hall effect in FMs & bilayers
- Detection of orbital accumulation in thin layers

Thanks to:



Leandro Salemi



Marco Berritta

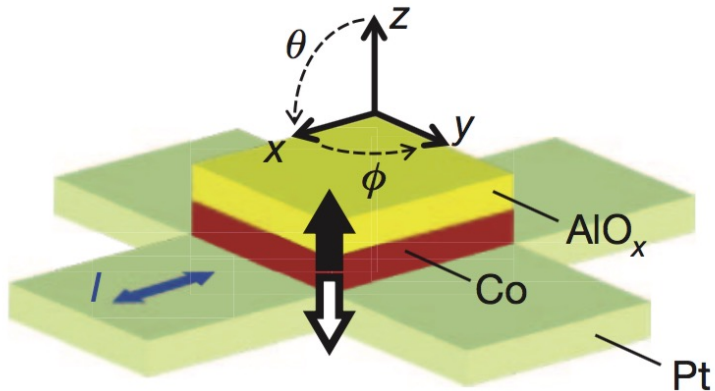


Sanaz Alikhah

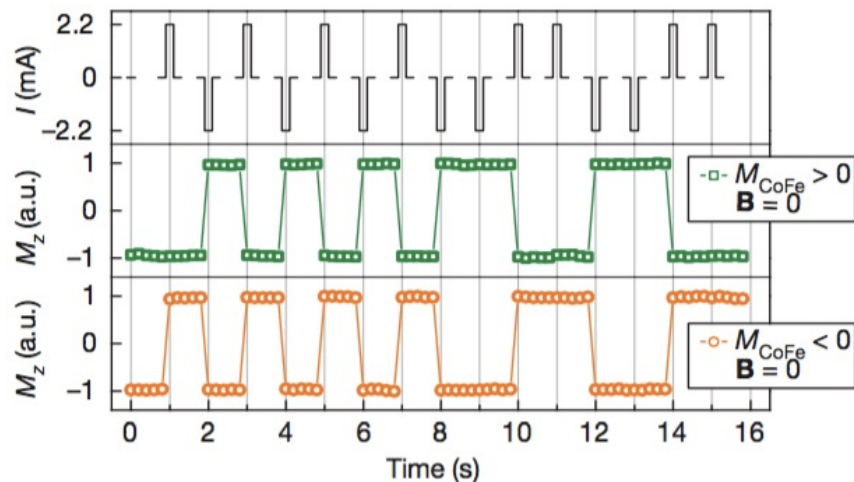
*Knut och Alice
Wallenbergs
Stiftelse*

Motivation - Spin-orbit torques

Magnetization switching with SOT



Spin Hall effect - small relativistic effect

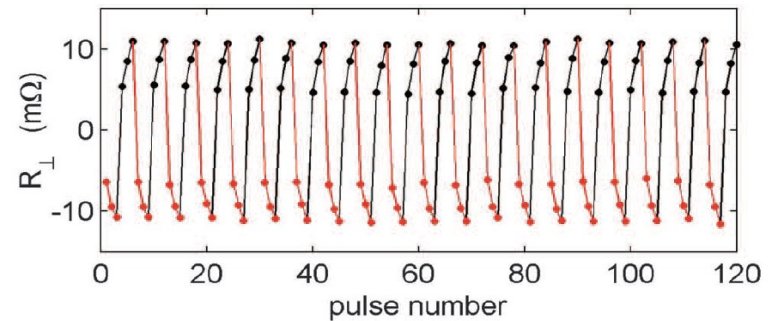


Miron et al, Nature **476**, 189 (2011)

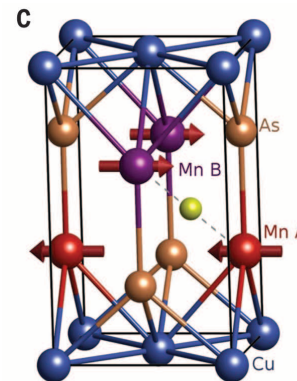
Liu et al, Science **336**, 555 (2012)

Review: Manchon et al, Rev.Mod.Phys. **91**, 035004 (2019)

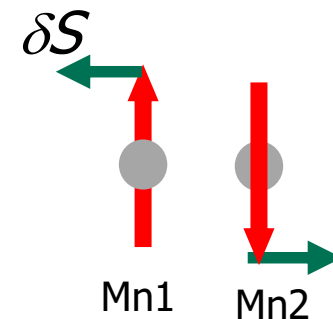
Staggered SOTs in special AFMs



CuMnAs



Rashba-Edelstein effect



Zelezny et al, PRL **113**, 157201 (2014)

Wadley et al, Science **351**, 587 (2016)

Zelezny et al, PRB **95**, 014403 (2017)

From a device view point

Spin polarized currents
Spin transfer torque (STT)



Spin currents
Spin-orbit torque (SOT)



Orbital currents
??

Slonczewski, 1996

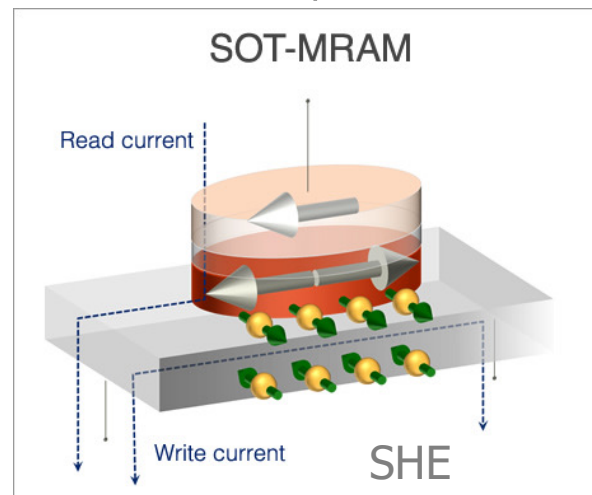
STT-MRAM



Everspin, Samsung, IBM

market-ready product
after ~20 years

Hirsch 1999 (SHE)
Miron et al, 2011

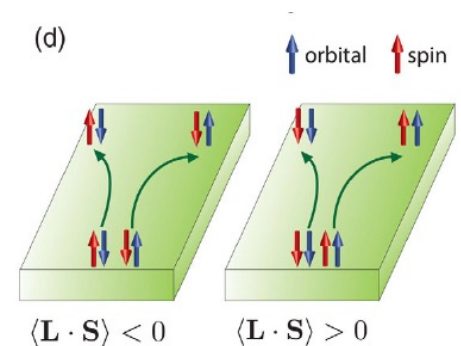


Technology being
developed for market

Observation SHE Pt:
Stamm et al, 2017

Go et al, ~2018 (OHE)

Orbital torque?

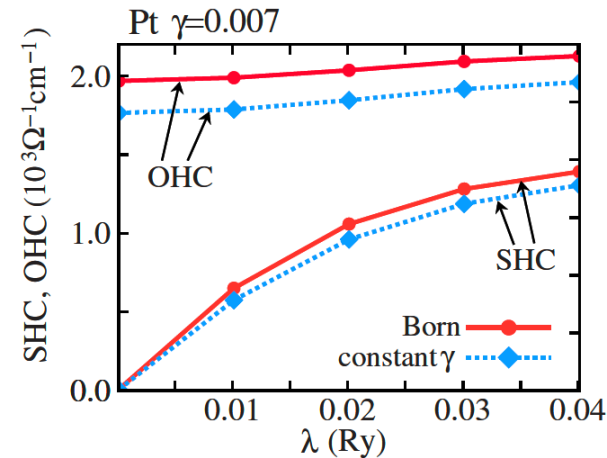


Current fundamental
research topic

Large orbital effects predicted

Orbital Hall effect predicted – without SOC

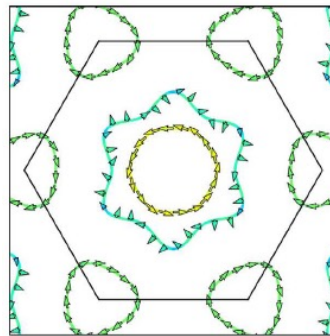
Tanaka et al, PRB **77**, 165117 (2008)



Orbital Rashba effect

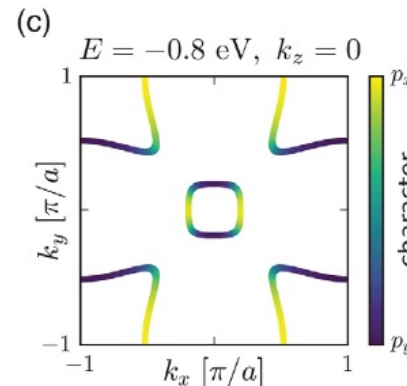
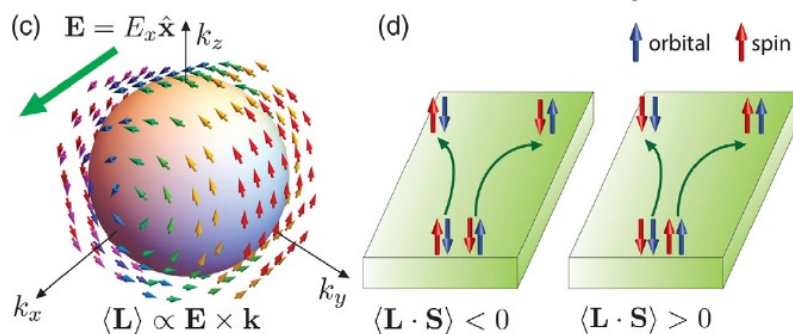
$$H_{\text{OR}}(\mathbf{k}) = \frac{\alpha_{\text{OR}}}{\hbar} \hat{\mathbf{L}} \cdot (\hat{\mathbf{z}} \times \mathbf{k})$$

Orbital Rashba Hamiltonian



Go, Lee, Mokrousov, Blügel et al, Sci. Rep. **7**, 46742 (2017)

Explanation of OHE

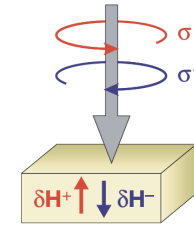


Orbital texture in k-space
Go, Jo, Kim & Lee, PRL **121**, 086602 (2018)

Electronic structure theory predictions

Orbital inverse Faraday effect –
not due to SOC

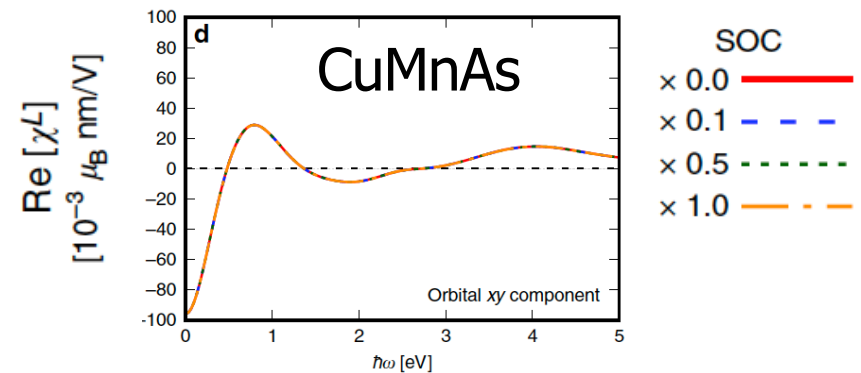
Berritta, Mondal, Carva, Oppeneer,
PRL **117**, 137203 (2016)



$$\vec{\delta M} \propto E^2 \vec{k}$$

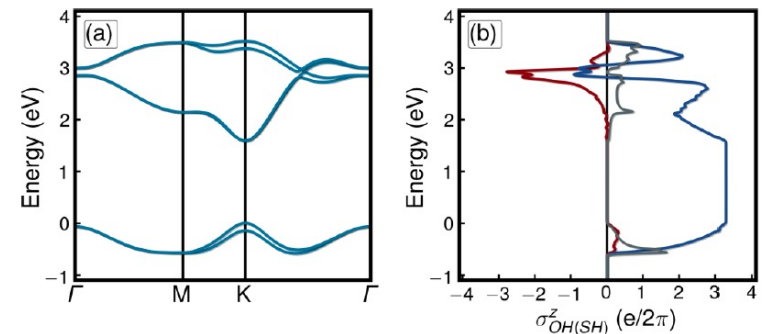
Orbital Rashba-Edelstein effect –
large, not due to SOC

Salemi, Beritta, Nandy, Oppeneer,
Nat. Commun. **10**, 5381 (2019)



Orbital Hall insul. phase in TMDCs

Canonico, Cysne, Molina, Muniz &
Rappoport, PRB **101**, 161409R (2020)



Many new papers in last years:

Go, Freimuth, Hanke, Xue et al, Phys.Rev.Res. **2**, 033401 (2020)

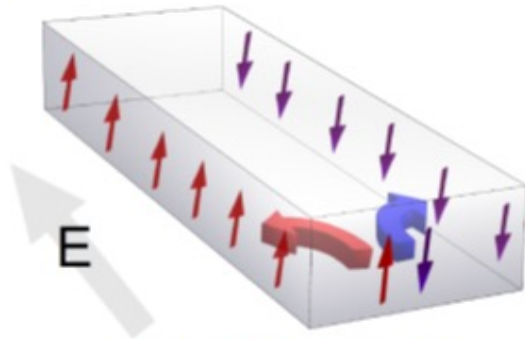
Go, Jo, Lee, Kläui & Mokrousov, EPL **135**, 37001 (2021)

Lee, Go, Park, Jeong et al, Nat. Commun. **12**, 6710 (2021)

Choi et al, arXiv 2109.14847

Phenomena: Charge-to-spin conversion – SHE & SREE

Spin Hall effect



Transport

$$J_y^{S_z} = \sigma_{yx}^{S_z} \cdot E_x$$

*Spin Conductivity
Tensor*

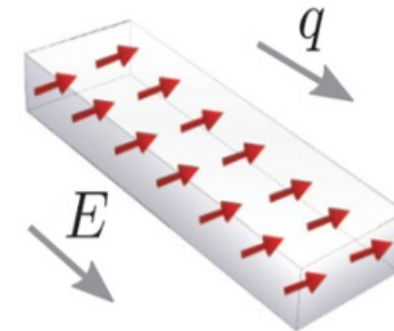
Dyakonov & Perel, JETP Lett. **13**, 467 (1971)

Hirsch, Phys.Rev.Lett. **83**, 1834 (1999)

Infinite bulk 3D crystal

<http://zfmezo.home.amu.edu.pl/research.php>

Rashba-Edelstein effect,
inv. spin-galvanic effect



Local

$$\delta S_y = \chi_{yx} \cdot E_x$$

*Spin Susceptibility
Tensor*

Edelstein, Solid State Comm. **73**, 233 (1990)

Originally: Rashba SOC + 2D symm.

Bychkov & Rashba, JETP Lett. **39**, 78 (1984)

$\sigma_{yx}^{S_z}$ & χ_{yx}
time-reversal even

Methodology – Electronic structure calculations

Rashba-Edelstein effect

(requires inv. symmetry breaking)

$$\delta S_i = \chi_{ij}^S E_j$$

Spin magneto-electric susceptibility

$$\delta L_i = \chi_{ij}^L E_j$$

Orbital magneto-electric susceptibility

DFT + linear response theory

$$\chi_{ij}^A = -\frac{ie}{m_e} \int_{\Omega} \frac{dk}{\Omega} \sum_{n \neq m} \frac{f_{nk} - f_{mk}}{\hbar\omega_{nmk}} \frac{A_{mnk}^i P_{nmk}^j}{-\omega_{nmk} + i\tau_{\text{inter}}^{-1}}$$

$A = \hat{L} \text{ or } \hat{S}$ (atomic sphere !)

Interband
(Fermi sea)

$$-\frac{ie}{m_e} \int_{\Omega} \frac{dk}{\Omega} \sum_n \frac{\partial f_{nk}}{\partial \epsilon} \frac{A_{nnk}^i P_{nnk}^j}{i\tau_{\text{intra}}^{-1}}$$

Intraband
(Fermi Surface)

$$\hbar\omega_{nmk} = \epsilon_{nk} - \epsilon_{mk}$$

Relativistic WIEN2k

(own implementation)

$$\delta \mathbf{M} = \mu_B \delta(\mathbf{L} + 2\mathbf{S})$$



Ab initio calculations

Spin Hall effect $\mathbf{J}^{S_k} = \boldsymbol{\sigma}^{S_k} \mathbf{E}$

$$A = \hat{\mathbf{J}}^{S_k} = \frac{\{\hat{S}_k, \hat{\mathbf{p}}\}}{2m_e V}$$

Orbital Hall effect $\mathbf{J}^{L_k} = \boldsymbol{\sigma}^{L_k} \mathbf{E}$

Spin/orbital current operator

Include scattering effects in an average through lifetimes $\tau_{intra}, \tau_{inter}$

(no explicit extrinsic effects such as side step or skew scattering)

SHE: Guo et al, PRL **100**, 096401 (2008)

OHE: Tanaka et al, PRB **77**, 165117 (2008)

Jo, Go, and Lee, PRB **98**, 214405 (2018)

Spin and orbital Nernst effects

$$\Lambda_{ij}^{S_k(L_k)} = \frac{\pi^2 k_B^2 T}{-3e} \left(\frac{d}{dE} \sigma_{ij}^{S_k(L_k)}(E) \right)_{E=E_F}$$

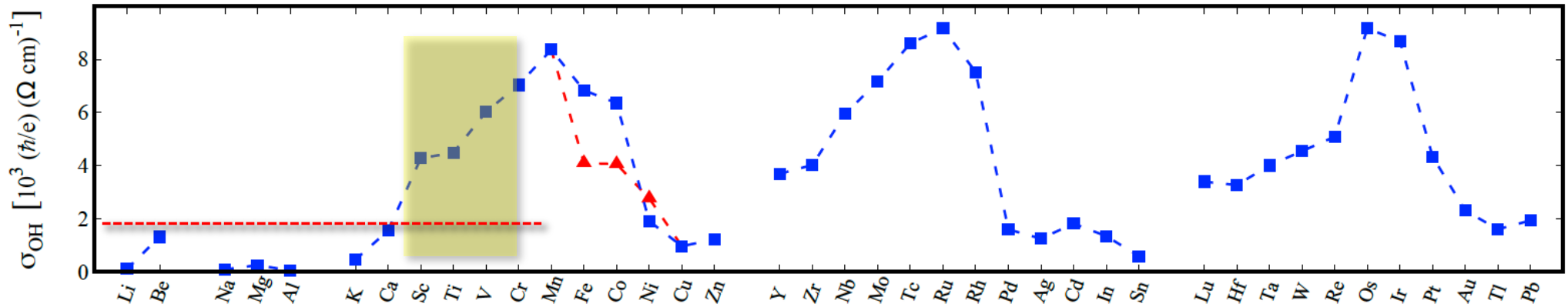
$$J_i^{S_k} = \sigma_{ij}^{S_k} E_j - \Lambda_{ij}^{S_k} \frac{dT}{dr_j}$$

Thermal gradient

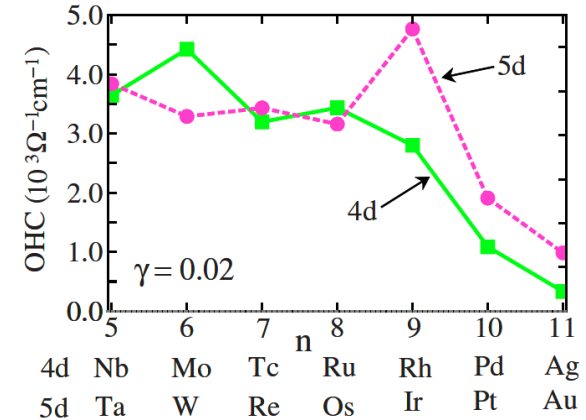
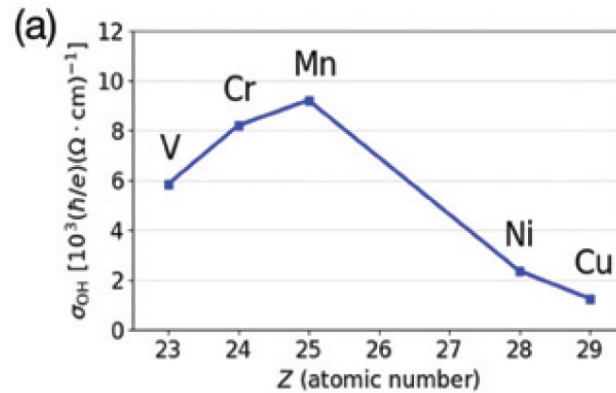
Spin Nernst effect: Meyer et al, Nat. Phys. **16**, 977 (2017)



Size of OHE for 40 elements



Jo et al, PRB **98**,
214405 (2018)



Tanaka et al, PRB **77**,
165117 (2008)

- Large effects for light 3d metals
- Good agreement with Jo et al (2018)
- Factor 2 & different trend from Tanaka et al (2008)
- Cheap, light metals for future orbitronics?

Recently predicted new effects

Hall

SHE	<u>OHE</u>
MSHE	MOHE

Rashba-Edelstein

SREE	<u>OREE</u>
MSREE	MOREE

Nernst

SNE	<u>ONE</u>
MSNE	MONE

$\sigma_{yx}^{S_z}$, χ_{yx}^S & $\Lambda_{yx}^{S_z}$
time-reversal even

	SHE	MSHE	OHE	MOHE
Require SOC	YES	YES	NO	YES
Require magnetism (<i>T</i> -odd)	NO	YES	NO	YES
Intraband (Fermi surf.)	NO	YES	NO	YES

Tanaka et al, PRB **77**, 165117 (2008)
 Sinova et al, Rev. Mod. Phys. **87**, 1213 (2015)
 Kimata et al, Nature **565**, 627 (2019)
 Salemi et al, Nat. Commun. **10**, 5381 (2019)

Edelstein, Solid State Comm. **73**, 233 (1990)
 Meyer et al, Nat. Phys. **16**, 977 (2017)
 Salemi and Oppeneer, PRB **106**, 024410 (2022)



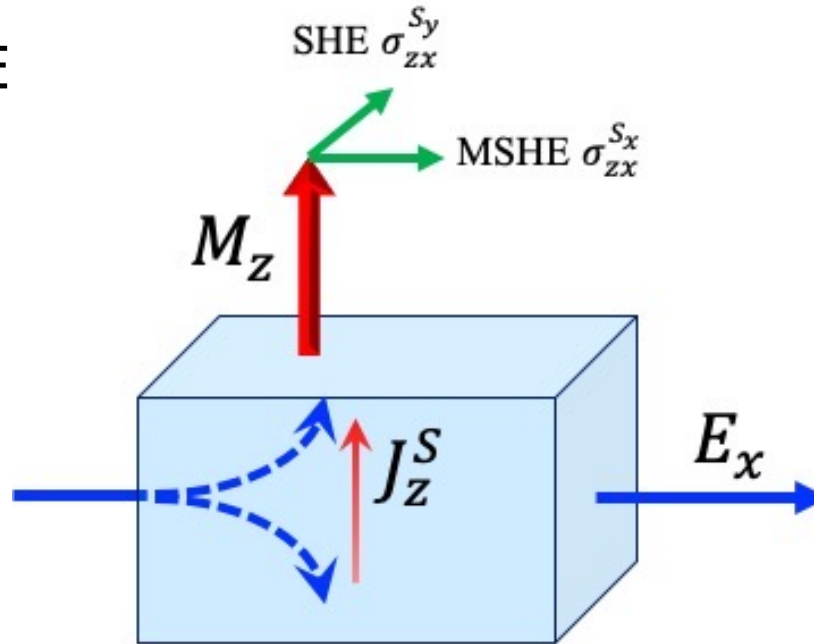
Unusual spin and orbital currents

Conventional SHE

$$\vec{J}^S \perp \vec{E} \perp \vec{S}$$

Time-reversal
and M even

Fermi sea or
interband



Spin Berry curvature

$$\sigma_{\text{OH(SH)}} = \frac{e}{\hbar} \sum_{n \neq m} \int \frac{d^3 k}{(2\pi)^3} (f_{m\mathbf{k}} - f_{n\mathbf{k}}) \Omega_{nm\mathbf{k}}^{X_z},$$

$$\Omega_{nm\mathbf{k}}^{X_z} = \hbar^2 \text{Im} \left(\frac{\langle u_{n\mathbf{k}} | j_y^{X_z} | u_{m\mathbf{k}} \rangle \langle u_{m\mathbf{k}} | v_x | u_{n\mathbf{k}} \rangle}{(E_{n\mathbf{k}} - E_{m\mathbf{k}} + i\eta)^2} \right)$$

Anomalous component

$$\vec{J}^S \perp \vec{E} \parallel \vec{S}$$

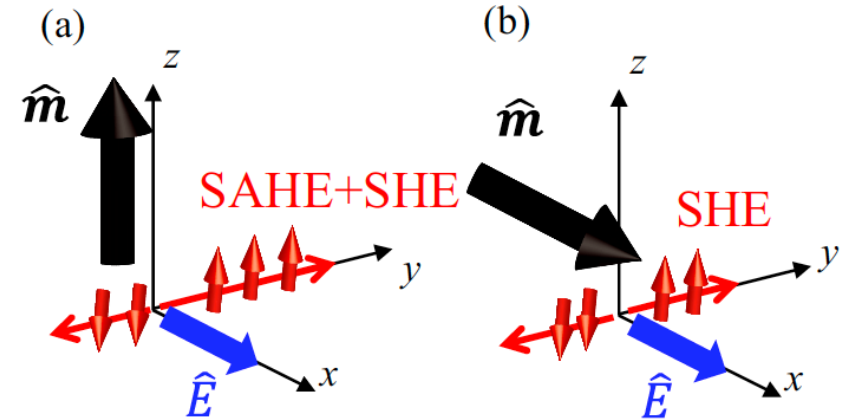
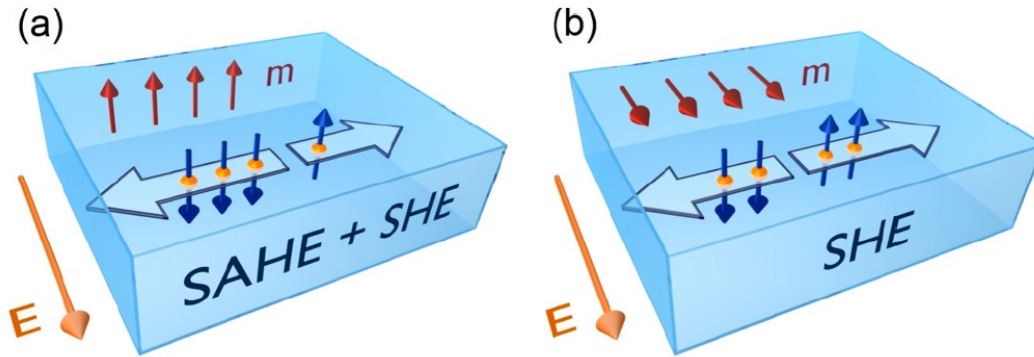
Time-reversal and
 M odd

Fermi surface or
intraband $n = m$

Calculations show:
Always present for
magnetic materials
(MSHE and MOHE)

Salemi, Berritta, Oppeneer,
PRMat. 5, 074407 (2021)

Spin anomalous Hall effect (SAHE) in ferromagnets



$$\vec{J}^S \perp \vec{E} \perp \vec{S}$$

$$\sigma_{\alpha\gamma}^{\beta} = -\frac{2e^2}{\hbar} \text{Im} \int \frac{d\mathbf{k}}{(2\pi)^3} \sum_{\substack{n \in \text{occ.} \\ m \in \text{unocc.}}} \frac{\langle \psi_m | \hat{Q}_{\alpha}^{\beta} | \psi_n \rangle \langle \psi_n | \hat{v}_{\gamma} | \psi_m \rangle}{(E_n - E_m)^2}$$

Amin, Li, Stiles, Haney, PRB **99**, 220405R (2019)

(Fermi sea or interband term)

$$\sigma_{yx}^{S_z}$$

SHE-like, T -even

Miura & Masuda, PRMat. **5**, L101402 (2021)

Dependence on M direction

Symmetry analysis

Nonmagnetic metal

$$J_i^{S_k} = \underbrace{\sigma_{ij}^{S_k}}_{\epsilon_{ijk} \sigma_{SH}} E_j$$

$$\sigma_{xy}^{S_z} = \sigma_{yz}^{S_x} = \sigma_{zx}^{S_y}$$

(Pt)

Ferromagnet

$$\sigma_{xy}^{S_z} \neq \sigma_{yz}^{S_x} \sim \sigma_{zx}^{S_y}$$

(Fe, Ni)

" $\sigma_{xy}^{SAHE}(\vec{m})$ "

$$\sigma_{xy}^{S_z} \neq \sigma_{yz}^{S_x} \neq \sigma_{zx}^{S_y}$$

(Co)

Amin, Li, Stiles, Haney, PRB **99**, 220405R (2019)
Miura & Masuda, PRMat. **5**, L101402 (2021)

$\mathbf{M} \parallel \mathbf{u}_z$

$$\sigma^{S_x} = \begin{pmatrix} 0 & 0 & \sigma_{xz}^{S_x} \\ 0 & 0 & \sigma_{yz}^{S_x} \\ \sigma_{zx}^{S_x} & \sigma_{zy}^{S_x} & 0 \end{pmatrix}$$

(and σ^{S_z})

$$\sigma^{S_y} = \begin{pmatrix} 0 & 0 & \sigma_{xz}^{S_y} \\ 0 & 0 & \sigma_{yz}^{S_y} \\ \sigma_{zx}^{S_y} & \sigma_{zy}^{S_y} & 0 \end{pmatrix}$$

SHE-like, T -even
 ϵ_{ijk}

MSHE-like, T -odd
Repeated indices

Calculated results Fe, Co, Ni

	T-even						T-odd			
	SHE			OHE			MSHE		MOHE	
	$\sigma_{yz}^{S_x}$	$\sigma_{zx}^{S_y}$	$\sigma_{xy}^{S_z}$	$\sigma_{yz}^{L_x}$	$\sigma_{zx}^{L_y}$	$\sigma_{xy}^{L_z}$	$\sigma_{xz}^{S_x}$	$\sigma_{zx}^{S_x}$	$\sigma_{xz}^{L_x}$	$\sigma_{zx}^{L_x}$
Fe	441	456	92	4697	4698	4707	-593	739	1343	848
Co	839	8	-44	5103	4718	4737	614	1074	-358	1356
Ni	1606	1543	824	3306	3297	3149	394	-290	-66	1033

$$\frac{\hbar}{e} (\Omega \text{ cm})^{-1}$$

$$\hbar\tau^{-1} = 40 \text{ meV}$$



SAHE

OHE >> SHE

SHE components strongly *M*-anisotropic (SAHE)

MSHE ~ SHE => must be taken into account

MOHE (new) exists, but smaller than OHE

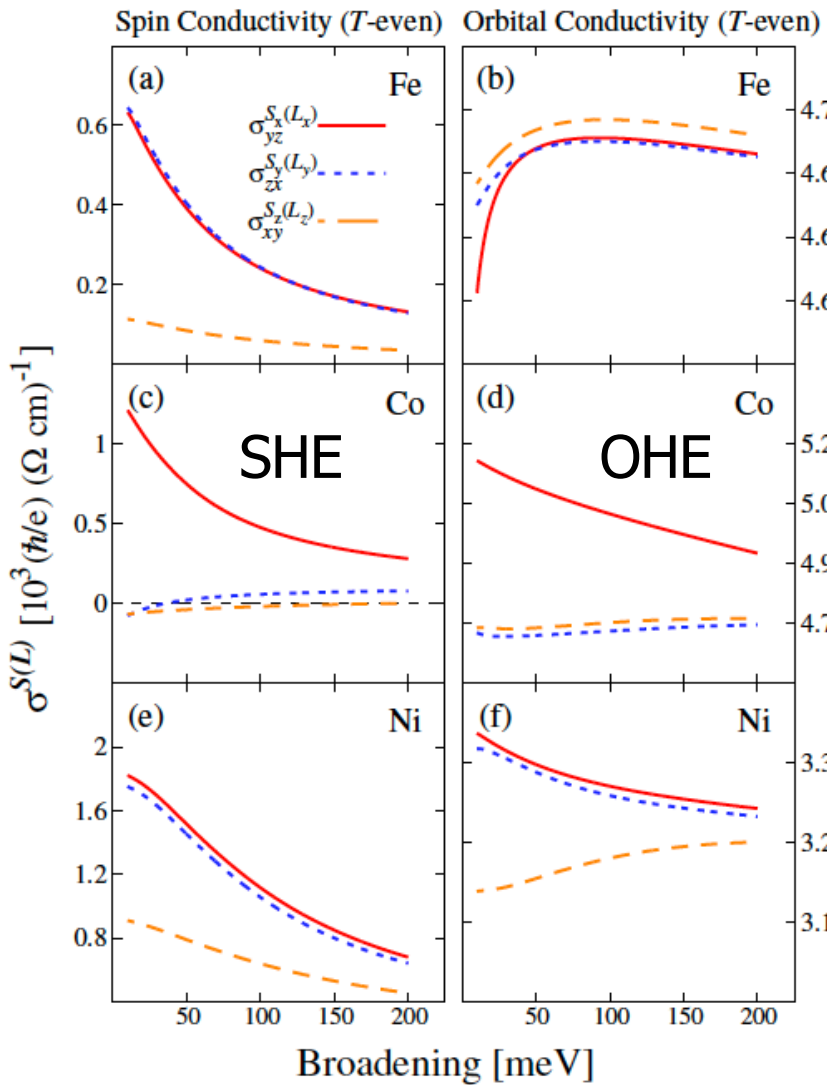
OHE present without SOC, all others require SOC

MSHE and MOHE (*intraband*) larger for pure samples

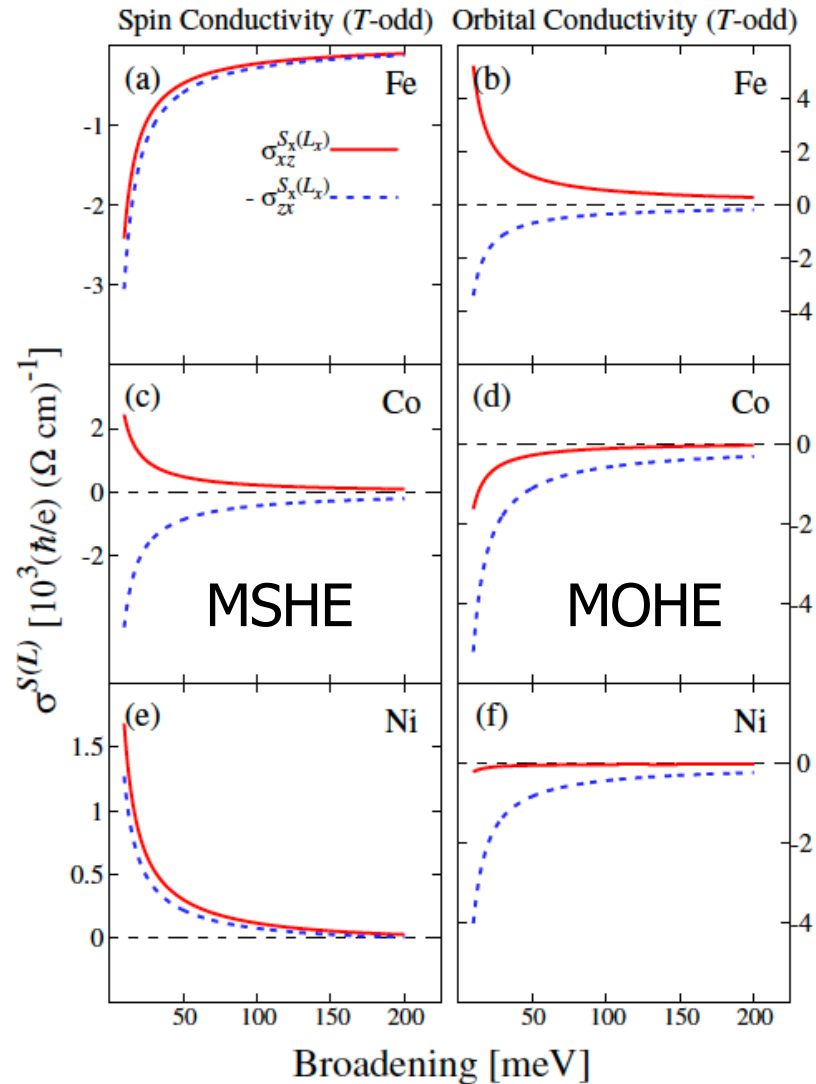


Lifetime dependence

T-even

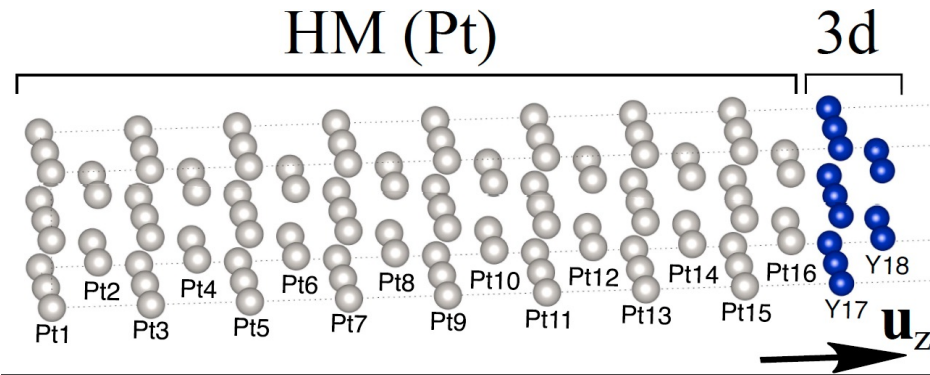


T-odd



MSHE and MOHE (*intraband*) much larger for pure samples

SOTs at symmetry-broken interface Pt/3d FM



$n \text{ Pt} / 2 \text{ Y}$
 $\text{Y} = \text{Ni, Co, Cu or Pt}$
 $n = 2 \text{ to } 16$

$$\delta \mathbf{S} = \chi^S \mathbf{E}$$

$$\mathbf{J}^{S\kappa} = \underline{\sigma^{S_k}} \mathbf{E}$$

\mathbf{E} along x

Effective torques

$$\mathcal{T} = \mathbf{M} \times \delta \mathbf{B} \quad \delta \mathbf{B} \approx |\mathbf{B}_{\text{XC}}| \frac{\delta \mathbf{S}}{|\mathbf{S}|}$$

$$\mathcal{T} = -2\mu_B |\mathbf{B}_{\text{XC}}| \left[\mathcal{M} \times \left(\chi^S \mathbf{E} \right) \right]$$

$$\mathcal{T}_{\text{FL}} \propto \mathbf{M} \times (\mathbf{E} \times \mathbf{u}_z)$$

$$\mathcal{T}_{\text{DL}} \propto \mathbf{M} \times [\mathbf{M} \times (\mathbf{E} \times \mathbf{u}_z)]$$

Ind. polarization

$$\left\{ \begin{array}{l} \delta S_y = \chi_{yx}^S E_x \quad T\text{-even} \\ \delta S_x = \chi_{xx}^S E_x \quad T\text{-odd} \end{array} \right.$$

$$\mathcal{M} = \mathbf{S}/|\mathbf{S}|$$

Time-rev. odd – field like

Time-rev. even – damping like

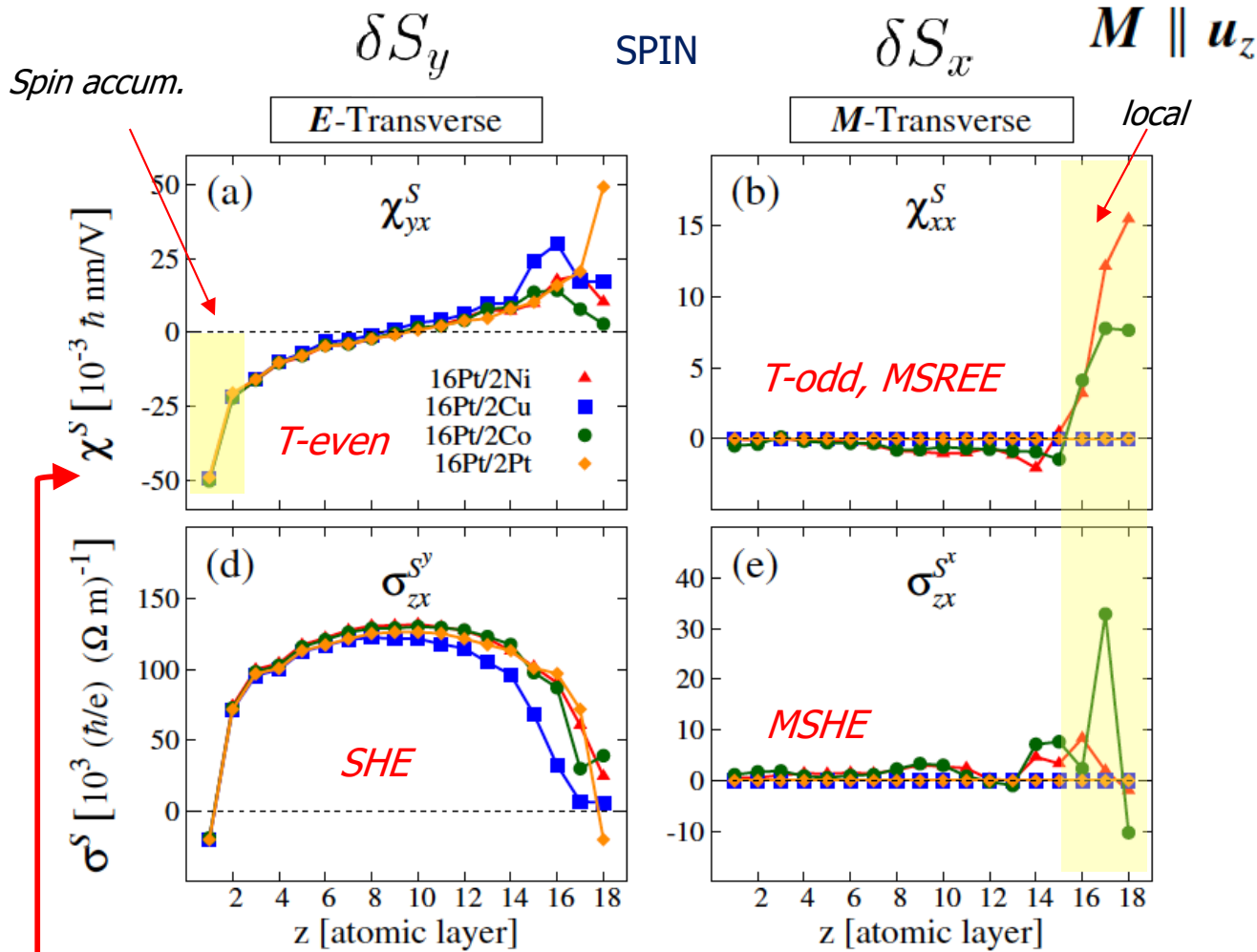
Early work:

Haney, Lee, Lee, Manchon, Stiles, PRB **88**, 214417 (2013)

Freimuth, Blügel, Mokrousov, PRB **90**, 174423 (2014)

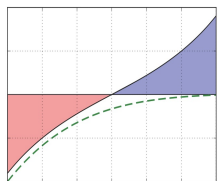


Layer-resolved induced spin polarization



- ❖ Typical transverse spin accumulation Pt
- ❖ Modified at the interface
- ❖ *M* and *T*-even effect

- ❖ Very local response at interface, along E_x
- ❖ Only exists for *magn.* material (*M* and *T*-odd)
- ❖ Same size as δS_y

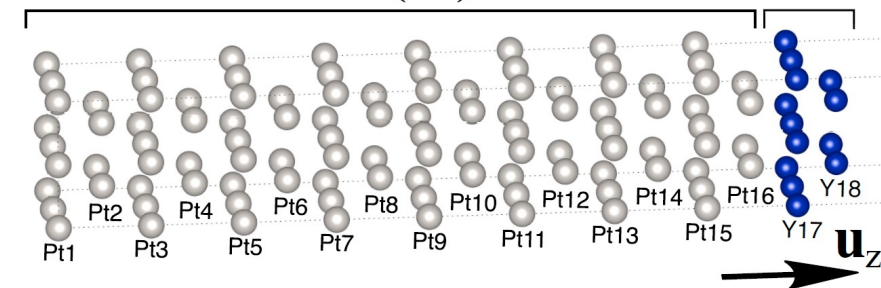


FL

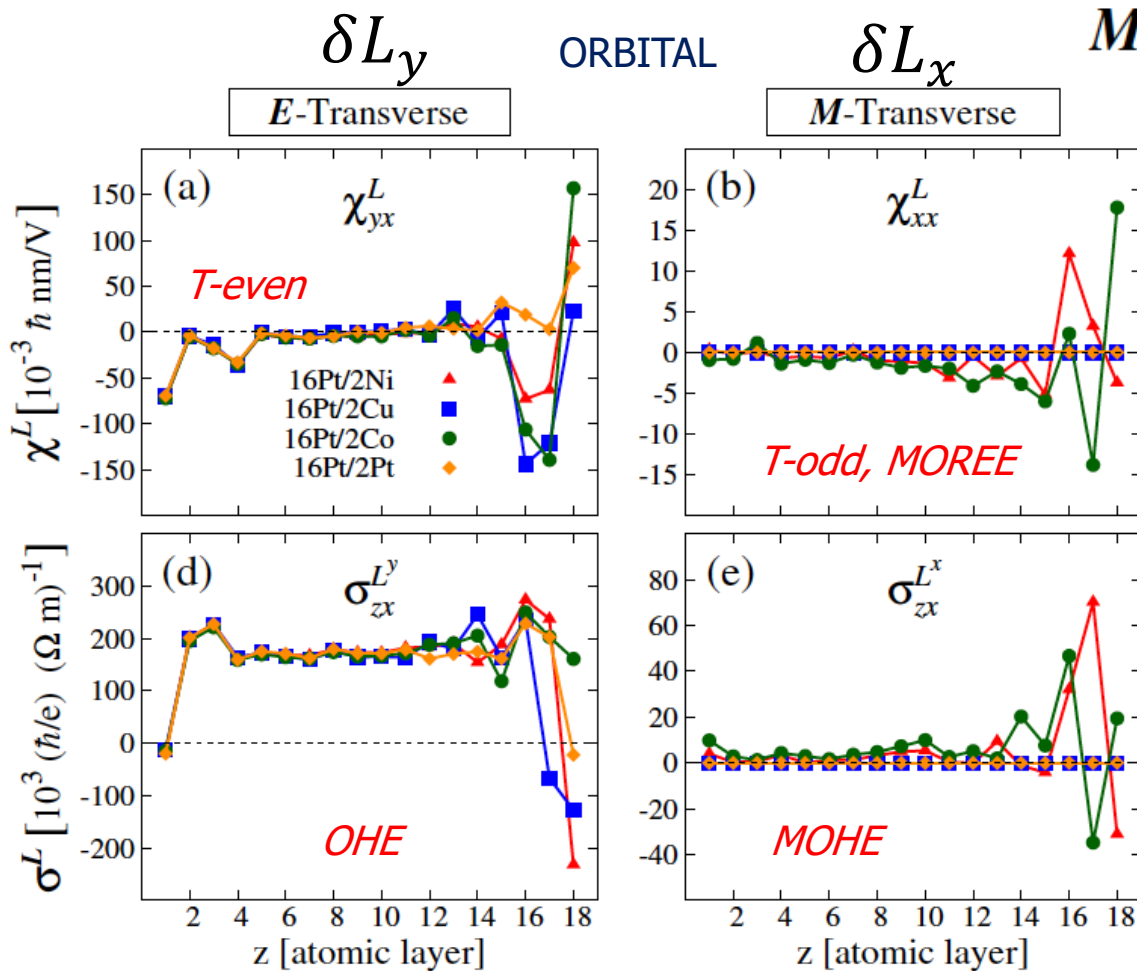
DL

HM (Pt)

3d



Results Pt/3d-bilayers –orbital polarization & current



- ❖ Huge OHE
- ❖ Orb. accumulation profile different from spin
- ❖ Enlarged at the interface
- ❖ M , T -even effect

- δL_x**
- ❖ Local response at interface, along E_x
 - ❖ Only exists for magn. material (M , T -odd)
 - ❖ Smaller than E -transv.

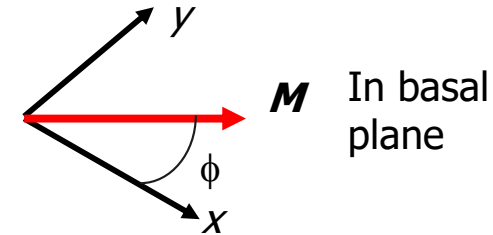
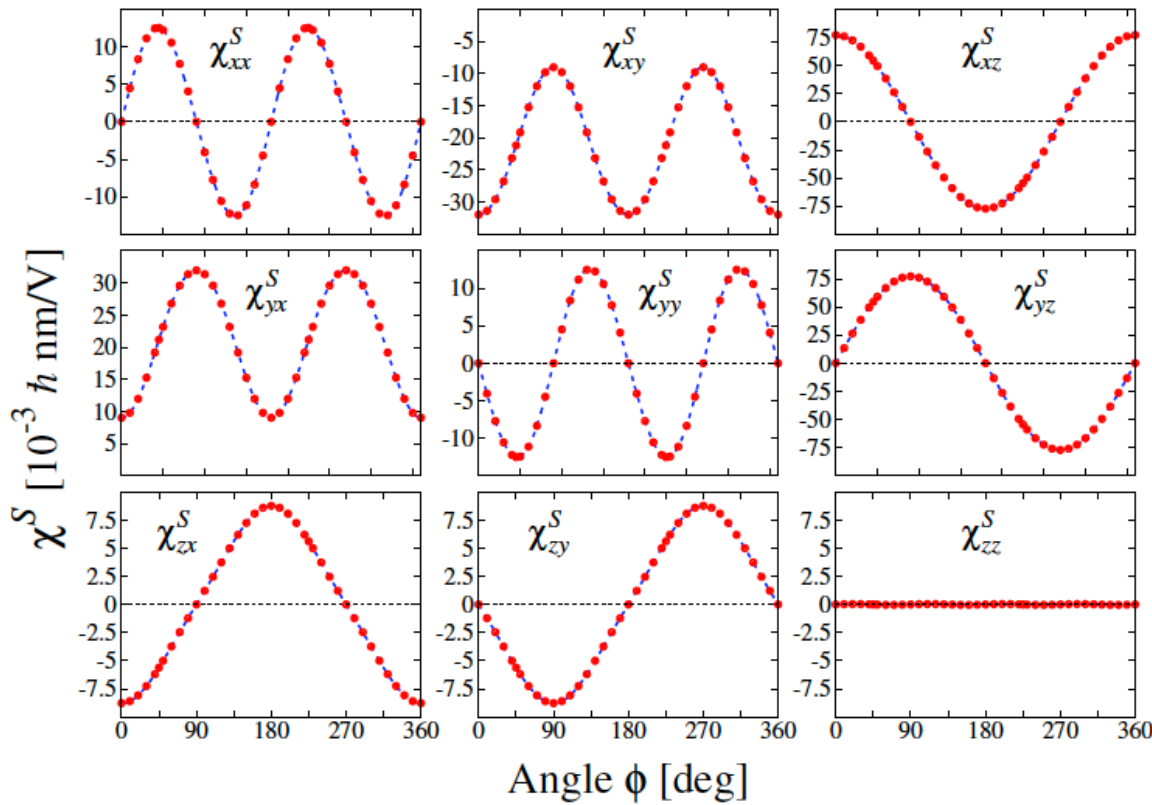
OHE and E -transv. orbital polarization not due to SOC



Magnetization direction dependence

Dependence of χ^S on M direction

12Pt/2Co, 1st Co



$$\mathcal{T} = -2\mu_B |B_{XC}| \left[\mathcal{M} \times \left(\chi^S E \right) \right]$$

*Induced spin polarization,
depends on direction of M*

Even in $M - 2\phi$ - FL torque
Odd in $M - \phi$ - DL torque

- Accurate angle dependence from *ab initio* calculations
- Identify even in M and odd in M components => analytical expressions
- Next step: full switching dynamics in the time domain

Many predicted effects – direct observations?

Hall

<u>SHE</u>	<u>OHE</u>
MSHE	MOHE

Rashba-Edelstein

SREE	<u>OREE</u>
MSREE	MOREE

Nernst

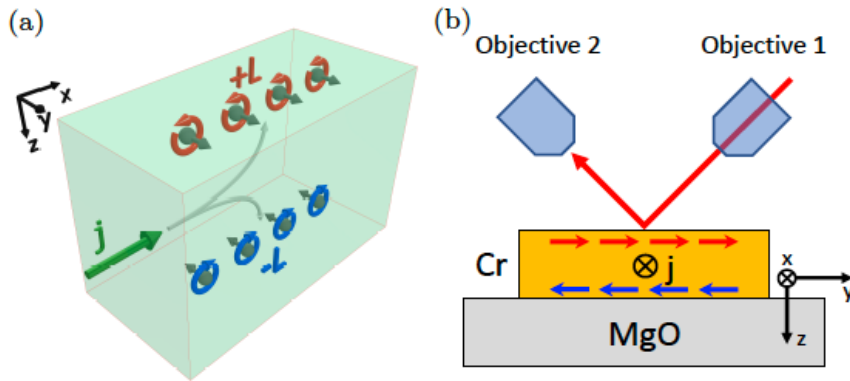
<u>SNE</u>	<u>ONE</u>
MSNE	MONE

Tanaka et al, PRB **77**, 165117 (2008)
 Sinova et al, Rev. Mod. Phys. **87**, 1213 (2015)
 Kimata et al, Nature **565**, 627 (2019)
 Salemi et al, Nat. Commun. **10**, 5381 (2019)

Stamm et al., PRL 119, 087203 (2017)
 Meyer et al, Nat. Phys. **16**, 977 (2017)
 Salemi and Oppeneer, PRB **106**, 024410 (2022)

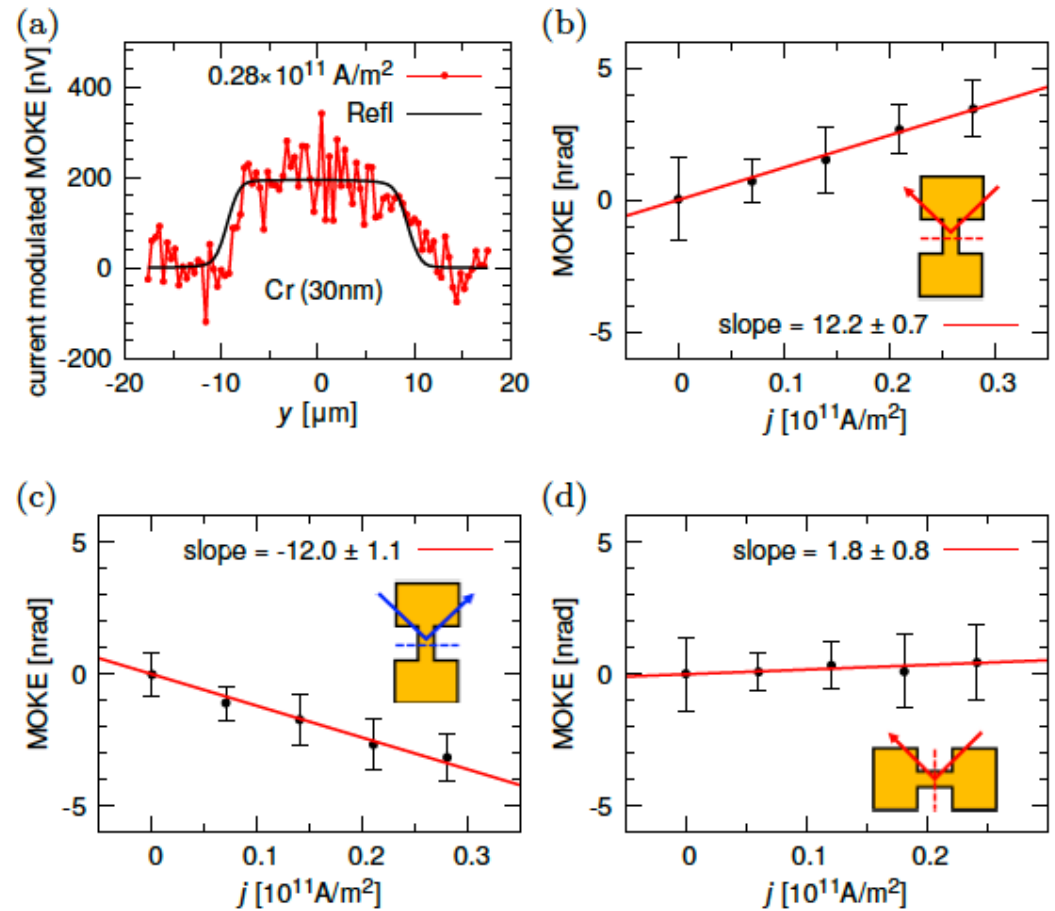
Experimental detection of orbital accumulation

L-MOKE on top surface



Cr film, $t = 30 \text{ nm}$, $w = 20 \text{ }\mu\text{m}$

Igor Lyalin, Roland Kawakami
(Columbus, Ohio)

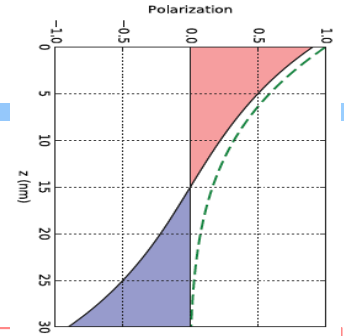


nanorads!

See also: Choi, Jo et al., Nature, in press (2023)



Electronic structure theory



OHE

Bulk complex MOKE

$$\theta_K^o = \frac{l_o \sigma_{zx}^{OH} \rho(t)^2 j D(E_F) e^{\frac{t}{2l_o}}}{\cosh(\frac{t}{2l_o})} \text{Re} \left\{ \Phi_K^{bulk,o} \kappa \left(\frac{(e^{-\kappa^- t} - 1) e^{-\frac{t}{l_o}}}{\kappa^-} - \frac{e^{-\kappa^+ t} - 1}{\kappa^+} \right) \right\}$$

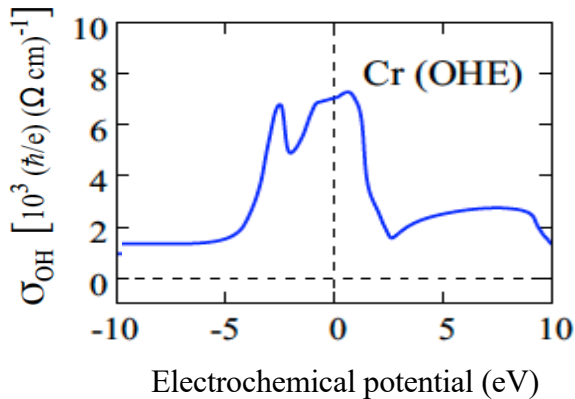
$$\kappa = (4\pi i \bar{n} \cos \psi) / \lambda$$

$$\kappa^\pm = \kappa \pm 1/l_o$$

$$\cos \psi = (1 - \sin^2 \phi_i / \bar{n}^2)^{1/2}$$

Orbital diffusion length
(parameter)

Refractive index



$$\sigma_{zx}^{OH} = 7000 \left(\frac{\hbar}{e}\right) [\Omega cm]^{-1}$$

$$\sigma_{zx}^{SH} \approx -70 \left(\frac{\hbar}{e}\right) [\Omega cm]^{-1} \quad \text{bcc Cr}$$

Stamm et al., PRL **119**, 087203 (2017)



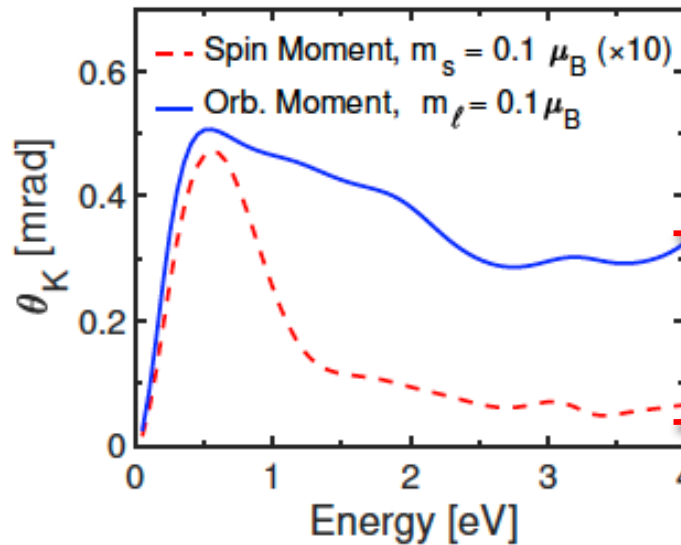
Ab initio based theory

Refractive index n from optical conductivity

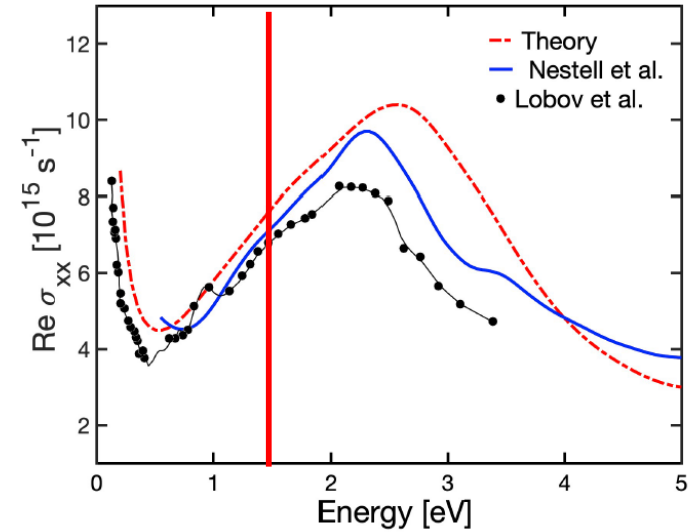
Bulk complex MOKE, using rel. DFT plus Zeeman field

$$\mathcal{H} = \mu_B \mathbf{B}_{ext} \cdot (\ell + 2s)$$

(b)



x 30
larger



$$\theta_K^L + i\epsilon_K^L = \frac{iQ\bar{n}n_0 \cos \phi_i \tan \phi_t}{\bar{n}^2 - n_0^2 \cos(\phi_i - \phi_t)}$$

WIEN2k FLAPW

$$\sigma_{\alpha\beta}(\omega) = -\frac{i\hbar}{V} \sum_{\mathbf{k}} \sum_{n \neq n'} \frac{f(\epsilon_{n\mathbf{k}}) - f(\epsilon_{n'\mathbf{k}})}{\epsilon_{n\mathbf{k}} - \epsilon_{n'\mathbf{k}}} \frac{\langle n'\mathbf{k} | j^\alpha | n\mathbf{k} \rangle \langle n\mathbf{k} | j^\beta | n'\mathbf{k} \rangle}{\hbar\omega - \epsilon_{n\mathbf{k}} + \epsilon_{n'\mathbf{k}} + i\hbar/\tau} + \sigma_D(\omega)\delta_{\alpha\beta}$$

Determination of orbital diffusion length in Cr

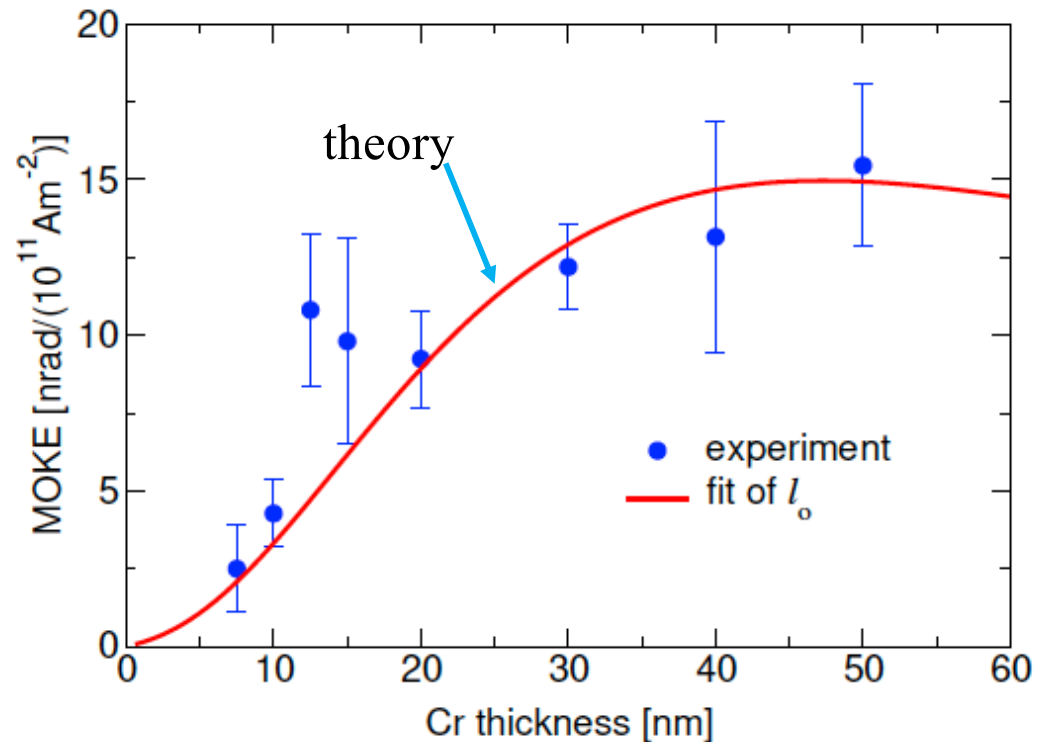
Fit MOKE or MOKE/ ρ^2 to obtain l_o

$$l_o = 6.6 \pm 0.6 \text{ nm}$$

$$\sigma_{zx}^{OH} = 7000 \left(\frac{\hbar}{e}\right) [\Omega \text{cm}]^{-1}$$

$$\theta_K^o = \frac{l_o \sigma_{zx}^{OH} \rho(t)^2 j D(E_F) e^{\frac{t}{2l_o}}}{\cosh\left(\frac{t}{2l_o}\right)} \times$$

$$\text{Re} \left\{ \Phi_K^{bulk,o} \kappa \left(\frac{(e^{-\kappa^- t} - 1)e^{-\frac{t}{l_o}}}{\kappa^-} - \frac{e^{-\kappa^+ t} - 1}{\kappa^+} \right) \right\}$$



Igor Lyalin, Roland Kawakami
(Columbus, Ohio)

Choi, Jo et al., Nature, in press (2023): Orbital diffusion length $l_o \sim 60 - 70$ nm Ti

But had to scale σ^{OH} by a factor of 1/100



Summary

Electronic structure calculations very powerful, predictive tool for development of spin-orbitronics and orbitronics

Ferromagnetic Fe, Co, Ni:

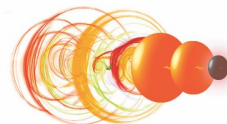
- Large magnetic spin Hall effect and MOHE coefficients (T -odd) predicted
- Expected to be present in other magnetic materials
- Leads to (novel) magnetic SNE and magnetic ONE (MSNE and MONE)

Pt/3d bilayers

- Presence of T -even and T -odd spin and orbital accumulations => diff. torques

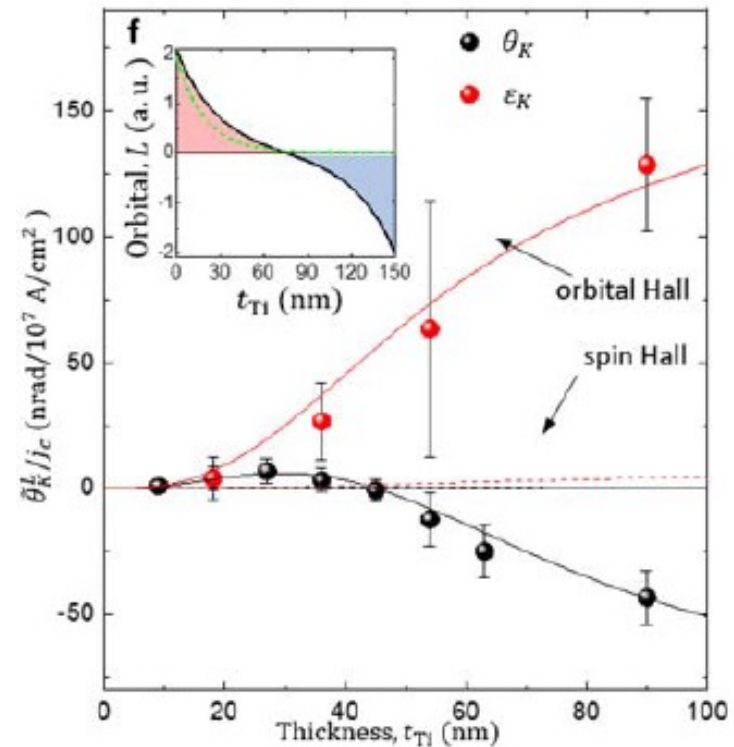
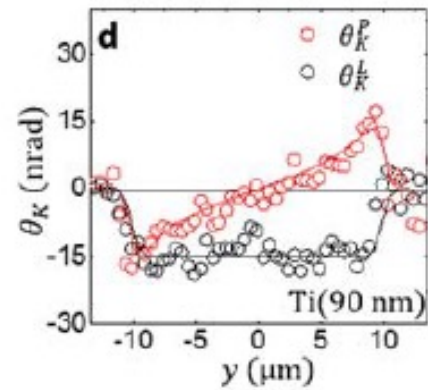
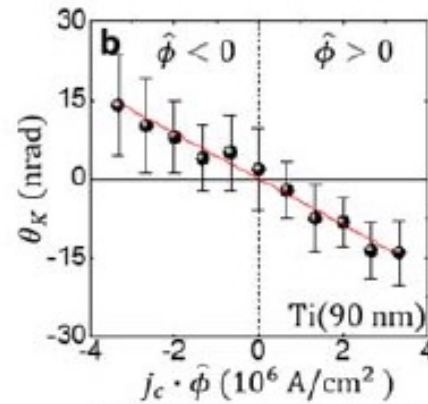
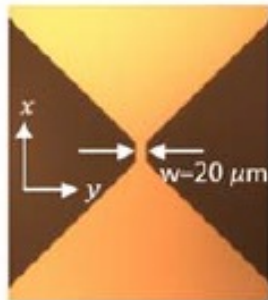
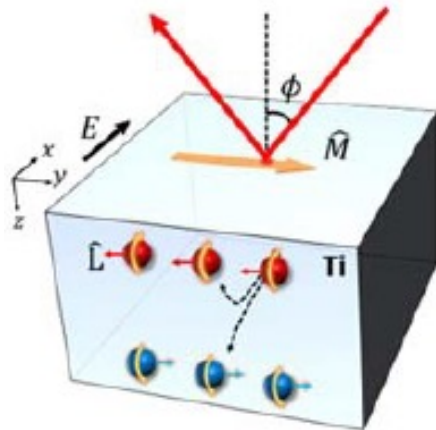
Elements

- Large OHE in light 3d metals predicted
- First evidence of very large OHE in Cr



Magneto-optical detection of OHE in Ti

L-MOKE

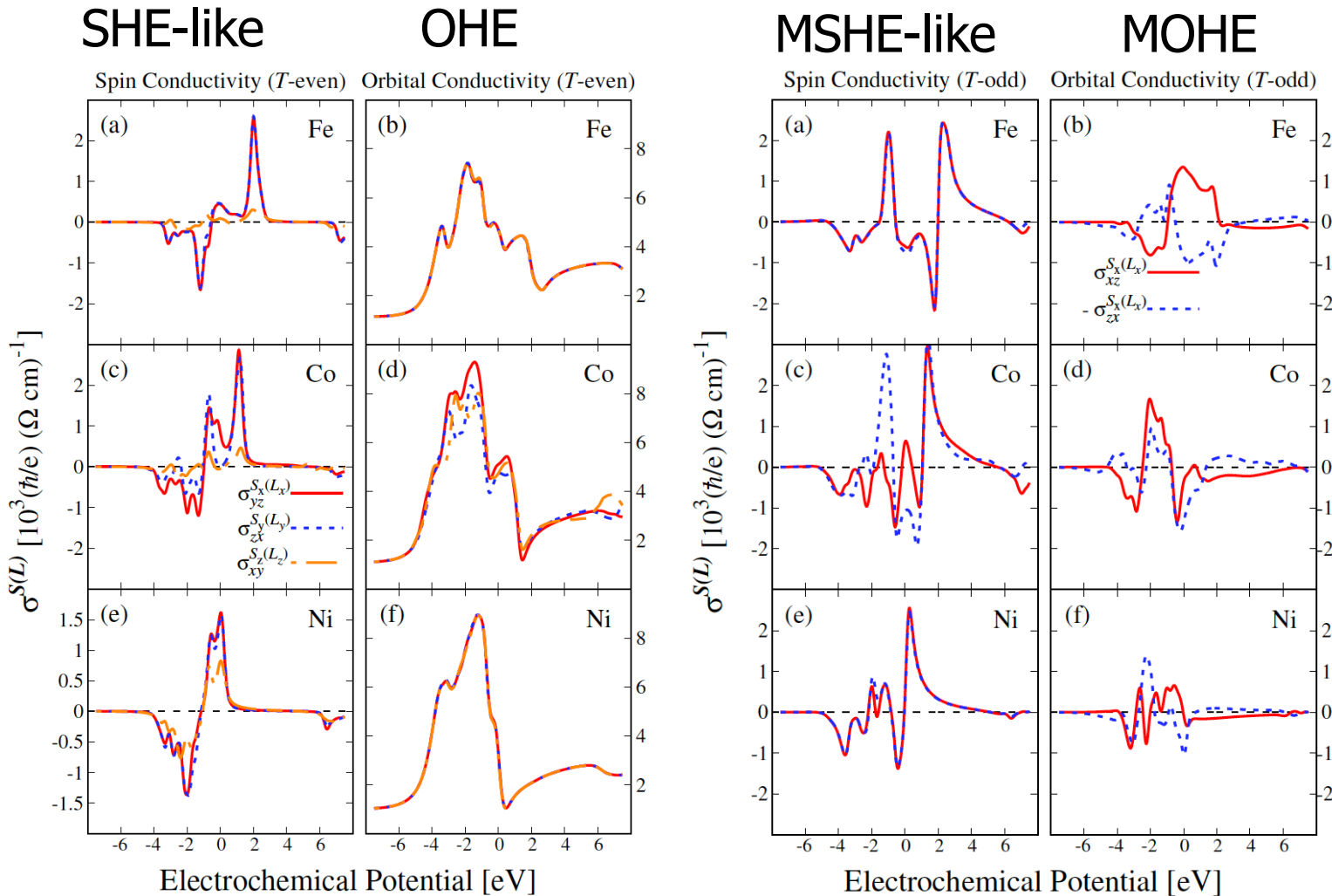


Orbital diffusion length $l_o \sim 60 - 70$ nm

But had to scale σ^{OH} by a factor of 1/100

Choi, Jo et al., Nature, in press (2023)

Ab initio calculated results Fe, Co, Ni



lifetime

$$\hbar\tau^{-1} = 40 \text{ meV}$$

OHE \ggg SHE
SHE *M*-anisotropic

MOHE \sim MSHE
MSHE \sim isotropic

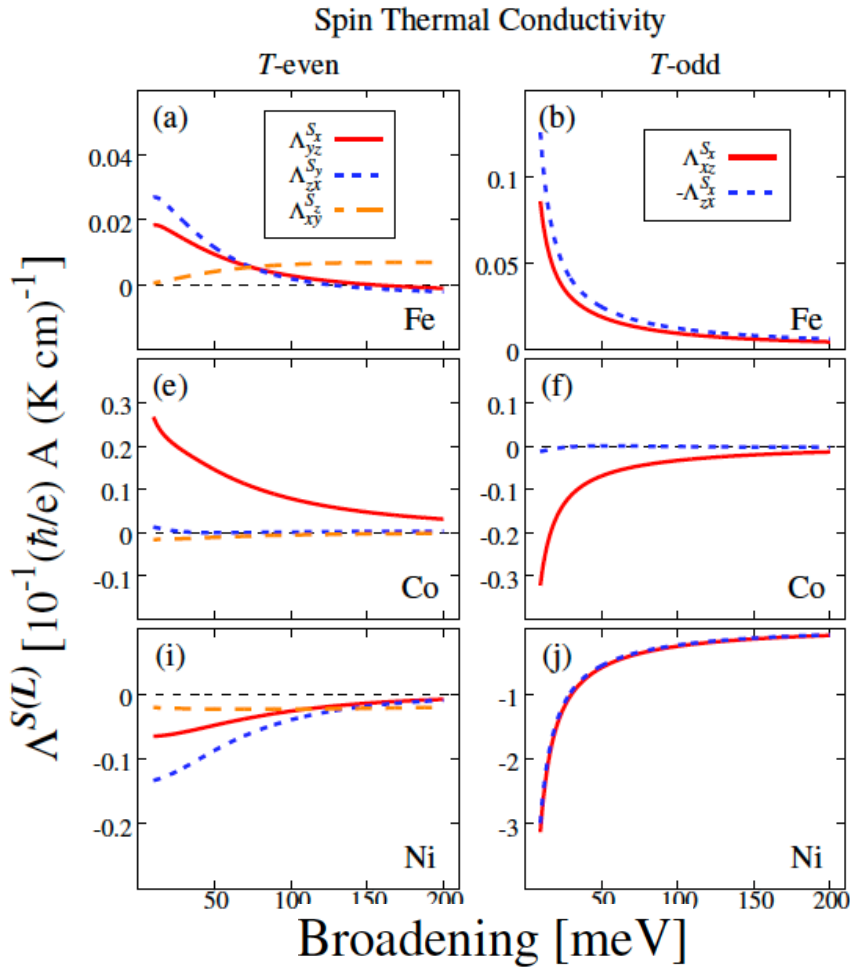
MSHE \sim SHE
MOHE $<$ OHE

Magnetic spin and orbital Nernst effects

$$J_i^{S_k} = \sigma_{ij}^{S_k} E_j - \Lambda_{ij}^{S_k} \frac{dT}{dr_j}$$

$$\Lambda_{ij}^{S_k(L_k)} = \frac{\pi^2 k_B^2 T}{-3e} \left(\frac{d}{dE} \sigma_{ij}^{S_k(L_k)}(E) \right)_{E=E_F}$$

$$\alpha_{ij}^{S_k(L_k)} = \frac{\Lambda_{ij}^{S_k(L_k)}}{\sigma_{ij}^{S_k(L_k)}}$$



“interband”

“intraband”

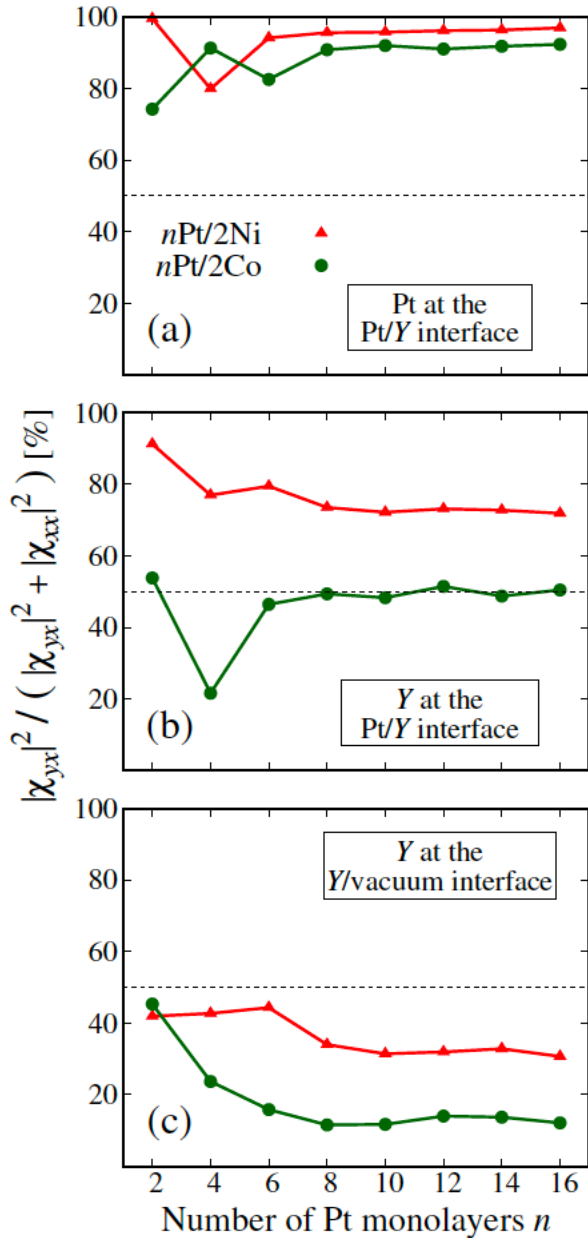
μVK^{-1}	SNC			MSNC	
	$\alpha_{yz}^{S_x}$	$\alpha_{zx}^{S_y}$	$\alpha_{xy}^{S_z}$	$\alpha_{xz}^{S_x}$	$\alpha_{zx}^{S_x}$
Fe	2.63	3.32	3.57	-3.94	-4.18
Co	19.72	1.10	28.74	-14.12	2.91
Ni	-3.35	-6.50	-2.80	-192.55	-248.60

$T = 300\text{ K}$

	ONC			MONC	
	$\alpha_{yz}^{L_x}$	$\alpha_{zx}^{L_y}$	$\alpha_{xy}^{L_z}$	$\alpha_{xz}^{L_x}$	$\alpha_{zx}^{L_x}$
Fe	2.87	3.02	2.93	1.07	-6.10
Co	-0.77	-0.83	-1.75	33.48	9.67
Ni	18.22	18.85	19.26	-253.57	-5.61

➤ Large MSNC and MONC predicted T-odd

Sizes of torques & Influence of SOC

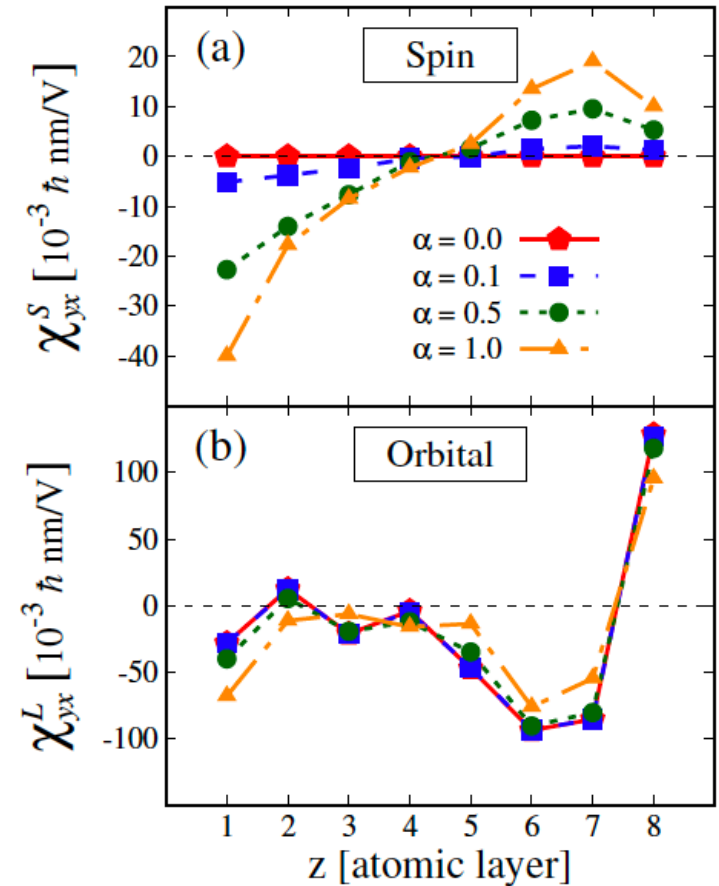


$$\frac{|\chi_{xy}^{S(L)}|^2}{|\chi_{xy}^{S(L)}|^2 + |\chi_{xx}^{S(L)}|^2}$$

Spin accum.

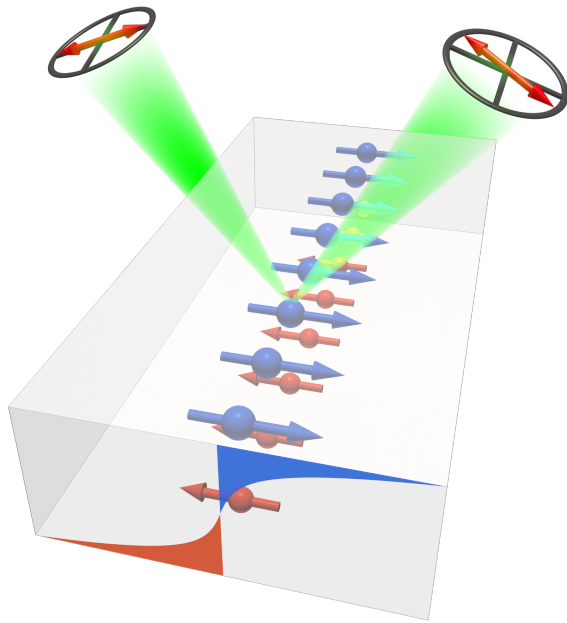
Magn. SHE

Influence of SOC

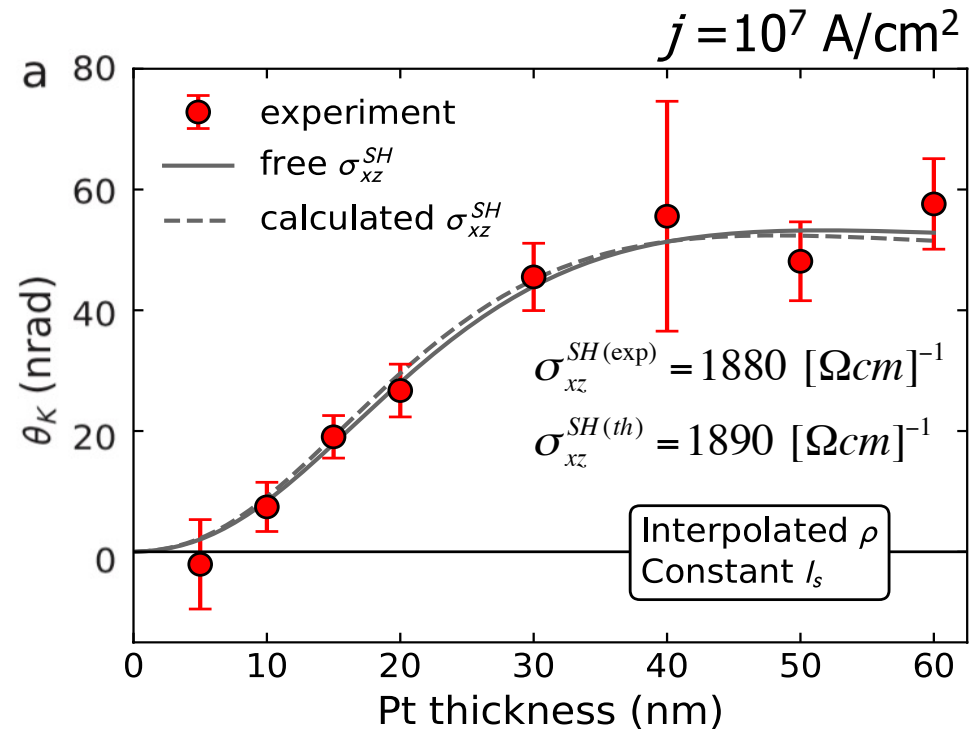


T -even orb. accumulation
and OHE not due to SOC,
all others SOC induced

Magneto-optical detection of SHE



Stamm, Murer, Berritta, Feng, Gabureac, Oppeneer & Gambardella, PRL **119**, 087203 (2017)



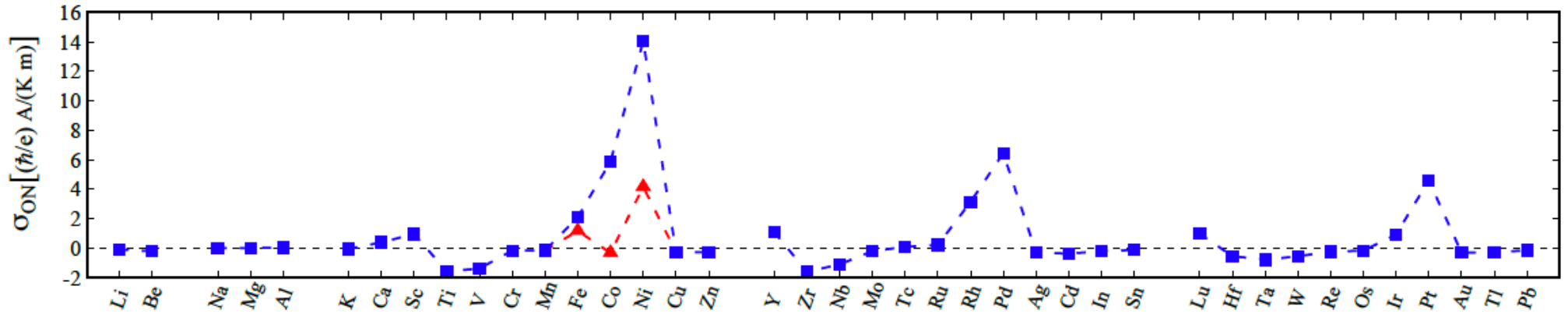
Excellent agreement with experiment!

Estimated $I_s = 11.4 \pm 2 \text{ nm}$ for pure Pt

- Accurate *ab initio* predictions of spin Hall effect possible
- Direct MOKE measurement of SH conductivity in heavy metals feasible



Orbital Nernst effect



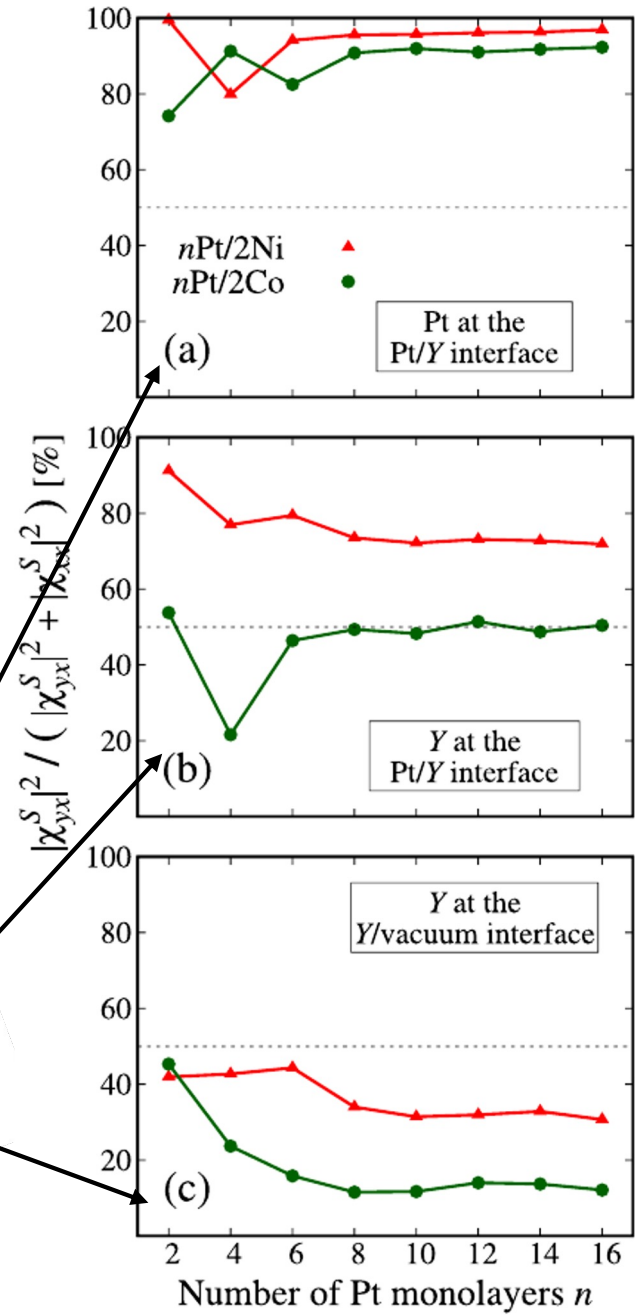
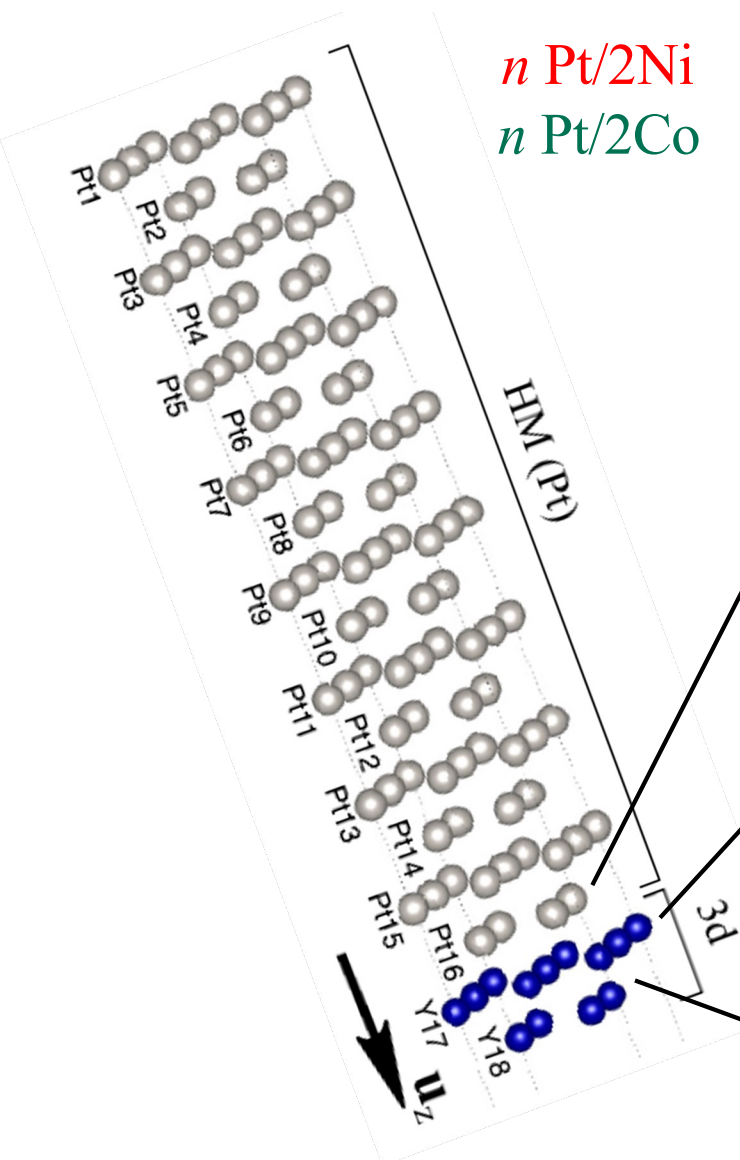
Large for nonmagnetic Ni, Pd & Pt

$$J_i^{S_k} = \sigma_{ij}^{S_k} E_j - \Lambda_{ij}^{S_k} \frac{dT}{dr_j}$$

Layer-resolved comparison of spin accumulation

$$\mathcal{T} = -2\mu_B |\mathbf{B}_{XC}| \left[\mathcal{M} \times (\chi^S \mathbf{E}) \right]$$

n Pt/2Ni
n Pt/2Co



$$|\chi_{yx}^S|^2 / (|\chi_{yx}^S|^2 + |\chi_{xx}^S|^2)$$

SHE dominated
 $\mathcal{T}_{\text{odd}} \gg \mathcal{T}_{\text{even}}$

$$\mathcal{T}_{\text{odd}} \gtrsim \mathcal{T}_{\text{even}}$$

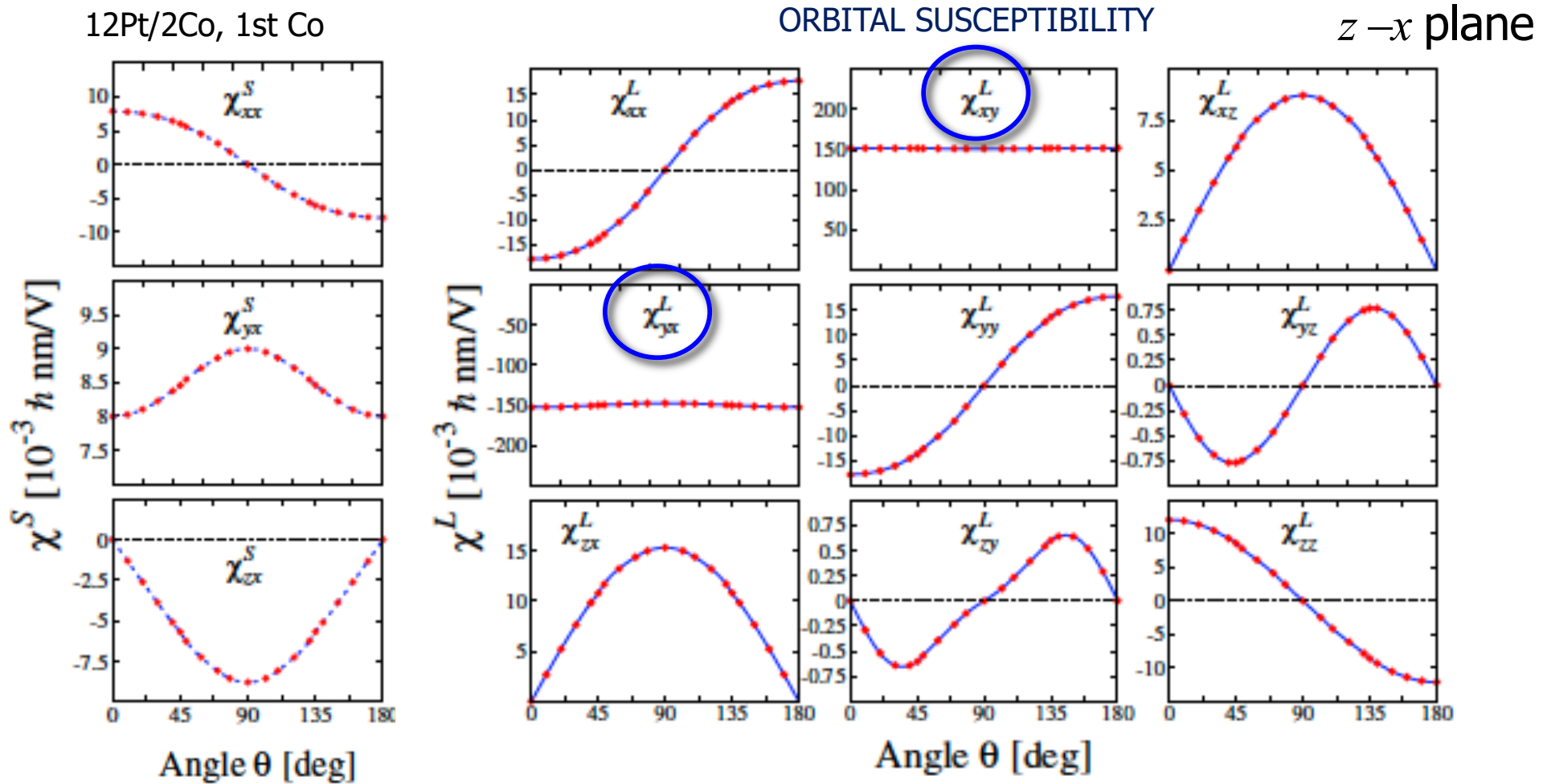
$$\mathcal{T}_{\text{odd}} \lesssim \mathcal{T}_{\text{even}}$$

Magn. SHE larger

But: lifetime dependent,
 χ_{xy} intraband



Induced orbital polarization – M direction dependence



- Orbital χ_{xy} is *nonrelativistic* and antisymmetric $\chi_{xy}^L \approx -\chi_{yx}^L$
- All other components due to SOC

Spin-charge angle tensor θ

$$\mathbf{J}^{S_k} = \sigma^{S_k} \sigma^{-1} \mathbf{J} = \underbrace{\sigma^{S_k} \rho}_{\text{red bracket}} \mathbf{J} = \frac{\hbar}{2e} \theta^{S_k} \mathbf{J} \quad \rho = \begin{pmatrix} \rho_1 & \rho_A & 0 \\ -\rho_A & \rho_1 & 0 \\ 0 & 0 & \rho_2 \end{pmatrix}$$

$\mathbf{M} \parallel \mathbf{u}_z$

$$\theta^{S_x} = \frac{2e}{\hbar} \begin{pmatrix} 0 & 0 & \sigma_{xz}^{S_x} \rho_2 \\ 0 & 0 & \sigma_{yz}^{S_x} \rho_2 \\ \sigma_{zx}^{S_x} \rho_1 - \sigma_{zy}^{S_x} \rho_A & \boxed{\sigma_{zx}^{S_x} \rho_A + \sigma_{zy}^{S_x} \rho_1} & 0 \end{pmatrix}$$

$$\theta^{S_y} = \frac{2e}{\hbar} \begin{pmatrix} 0 & 0 & \sigma_{xz}^{S_y} \rho_2 \\ 0 & 0 & \sigma_{yz}^{S_y} \rho_2 \\ \sigma_{zx}^{S_y} \rho_1 - \sigma_{zy}^{S_y} \rho_A & \sigma_{zx}^{S_y} \rho_A + \sigma_{zy}^{S_y} \rho_1 & 0 \end{pmatrix}$$

$$\theta^{S_z} = \frac{2e}{\hbar} \begin{pmatrix} \sigma_{xx}^{S_z} \rho_1 - \sigma_{xy}^{S_z} \rho_A & \boxed{\sigma_{xx}^{S_z} \rho_A + \sigma_{xy}^{S_z} \rho_1} & 0 \\ \sigma_{yx}^{S_z} \rho_1 - \sigma_{yy}^{S_z} \rho_A & \sigma_{yx}^{S_z} \rho_A + \sigma_{yy}^{S_z} \rho_1 & 0 \\ 0 & 0 & \sigma_{zz}^{S_z} \rho_2 \end{pmatrix}$$

Mixing of effects

Spin-filtering on AHE